

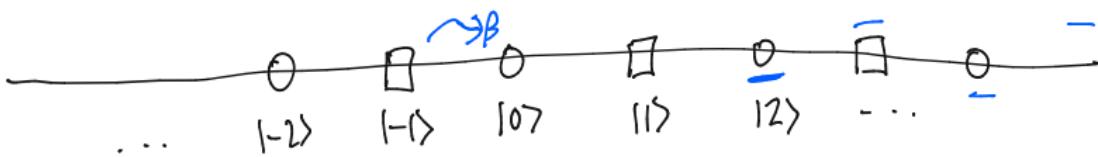
**PHYS 4410**  
**Quantum Mechanics 2**  
**Spring 2023**

**Lecture 12**

**Metals and insulators**

February 17

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$\circ$  and  $\square$  inequivalent:

$$H = \sum_{n=-\infty}^{\infty} \underbrace{[-\beta |n\rangle \langle n+1| - \beta |n+1\rangle \langle n|]}_{\text{local}} - (-1)^n \gamma |n\rangle \langle n| \quad ].$$

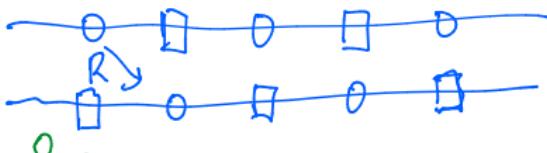
$$H|n\rangle = -\beta |n-1\rangle - \beta |n+1\rangle \begin{cases} -\gamma |n\rangle & n \text{ even} \\ +\gamma |n\rangle & n \text{ odd} \end{cases}$$

$$H = H_\beta + H_\gamma$$

$\uparrow$   $\uparrow$   
 $\beta$ -terms  $\gamma$ -term  
 (new for today)

2 Let  $R = \sum_{n=-\infty}^{\infty} |n\rangle\langle n|$ . Does  $[H, R] = 0$ ?  
 $[H, R] \neq 0$ .

Intuitively:



$$[H_B + H_\gamma, R] = \cancel{[H_B, R]} + [H_\gamma, R]$$

so  $[H_\gamma, R] \neq 0$ .

$$\begin{aligned} H_\gamma R |n\rangle &= H_\gamma |n+1\rangle = -\gamma(-1)^{n+1} |n+1\rangle \\ RH_\gamma |n\rangle &= R \left[ -\gamma \underbrace{(-1)^n}_{\text{underbrace}} |n\rangle \right] = -\gamma(-1)^n |n+1\rangle \end{aligned} \quad \left. \begin{array}{l} \text{differ by } -1. \\ \text{---} \end{array} \right\}$$

Fix:

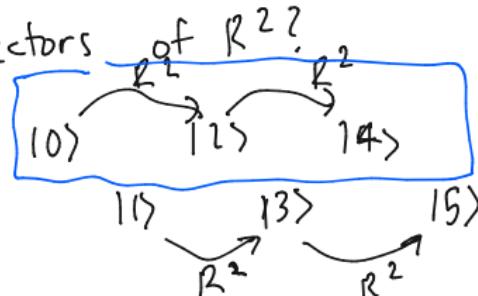
$$\begin{array}{c} \text{---} \\ \text{---} \end{array} \quad \begin{array}{ccccccc} \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} \\ \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} \end{array} \quad R \cdot R = R^2.$$

You can check:  $[H, R^2] = 0$ .

discrete translation sym:  $\{1, R^2, R^4, \dots, R^{-2}, R^{-4}, \dots\}$   
 subgroup of  $\{1, R, R^2, \dots\}$

3 What are eigenvalues/vectors of  $R^2$ ?

$$R^2 = \sum_{n=-\infty}^{\infty} |n+2\rangle\langle n|$$



Claim:  $|\tilde{k}\rangle_e = \sum_{n \text{ even}} e^{ikn} |n\rangle$

$\xrightarrow{\text{always orthogonal}}$

$|\tilde{k}\rangle_o = \sum_{n \text{ odd}} e^{ik(n-1)} |n\rangle$

is equivalent to  
 $(k \sim k + 2\pi)$

Have 2 eigenvectors for each  $k$ .

When  $[H, R] = 0$ , Brillouin zone: independent  $k$  cure

$$0 \leq k < 2\pi$$

$$-\pi \leq k < \pi$$

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If  $[H, R^2]$ :  $\sum_{n \text{ even}} e^{ikn} |n\rangle = \sum_{n \text{ even}} e^{i\ln(k+\pi)} |n\rangle$

$e^{2\pi i \times \text{integer}} = 1$

Now:  $-\pi/2 < k \leq \pi/2$ .

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$$\langle H | n \rangle = -\beta |n-1\rangle - \beta |n+1\rangle - \gamma (-1)^n |n\rangle$$

Since  $[H, R^2] = 0$ , know eigenvectors of  $H$  are eigenv. of  $R^2$ :

$$c_e |\bar{k}\rangle_e + c_o |\bar{k}\rangle_o$$

$$H [c_e |\bar{k}\rangle_e + c_o |\bar{k}\rangle_o] = H \sum_{m=-\infty}^{\infty} e^{2ikm} [c_e |2m\rangle + c_o |2m+1\rangle]$$

$$= \sum_{m=-\infty}^{\infty} e^{2ikm} \left[ -\beta c_e |2m-1\rangle - \beta c_e |2m+1\rangle - \gamma c_e |2m\rangle \right. \\ \left. - \beta c_o |2m\rangle - \beta c_o |2m+2\rangle + \gamma c_o |2m+1\rangle \right]$$

e.g.

$$\sum_{m=-\infty}^{\infty} e^{2ikm} (-\beta c_e) |2m-1\rangle \xrightarrow{m \geq m'+1} \sum_{m'=-\infty}^{\infty} e^{2ik} \cdot e^{2ikm'} (-\beta c_e) |2m'+1\rangle$$

$$= \sum_{m=-\infty}^{\infty} e^{2ikm} \left[ -\gamma c_e |2m\rangle + \gamma c_o |2m+1\rangle - \beta c_e (1 + e^{2ik}) |2m+1\rangle \right. \\ \left. - \beta c_o (1 + e^{-2ik}) |2m\rangle \right] \\ = (-\gamma c_e - \beta c_o (1 + e^{-2ik})) |\bar{k}\rangle_e + (\gamma c_o - \beta (1 + e^{2ik})) |\bar{k}\rangle_o.$$

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$$\underbrace{H|\psi\rangle = E|\psi\rangle}_{\downarrow}, \quad \text{where } |\psi\rangle = c_1|\bar{k}\rangle_e + c_2|k\rangle_o$$

$$\begin{aligned} E c_e &= -\gamma c_e - \beta(1+e^{-2ik})c_o \\ E c_o &= +\gamma c_o - \beta(1+e^{2ik})c_e \end{aligned}$$

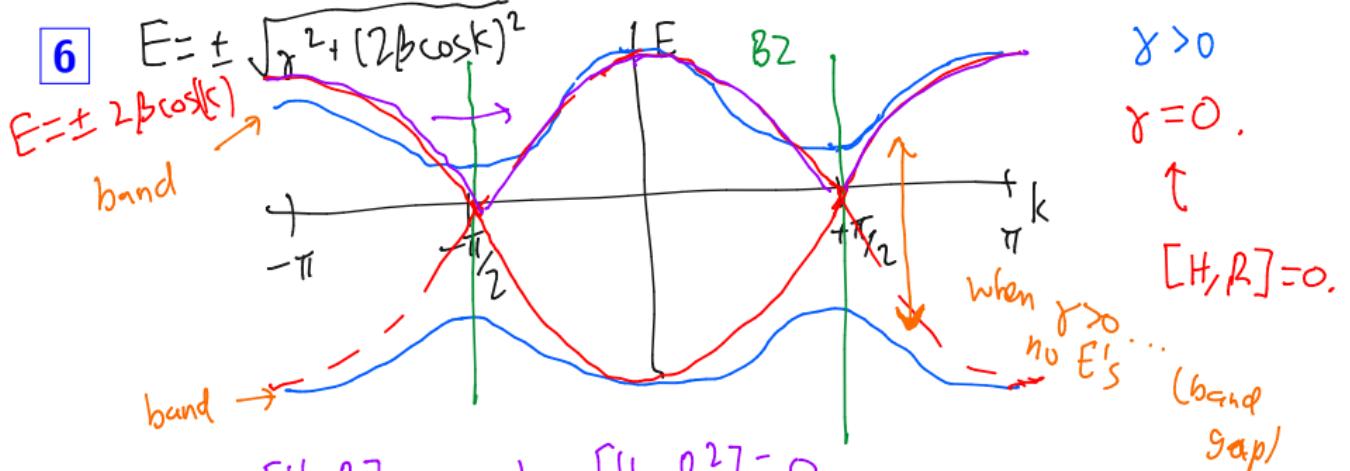
$$E \begin{pmatrix} c_e \\ c_o \end{pmatrix} = \underbrace{\begin{pmatrix} -\gamma & -\beta(1+e^{-2ik}) \\ -\beta(1+e^{2ik}) & \gamma \end{pmatrix}}_{Z\beta C} \begin{pmatrix} c_e \\ c_o \end{pmatrix}$$

$$\begin{aligned} \text{Find } E \text{ by: } \det(EI - Z) &= 0 = \det \begin{pmatrix} E+\gamma & Z\beta C \\ \beta C & E-\gamma \end{pmatrix} \\ &= (E+\gamma)(E-\gamma) - \beta^2(1+e^{-2ik})(1+e^{2ik}) \\ &= E^2 - \gamma^2 - \beta^2(2 + 2\cos(2k)) \end{aligned}$$

$$\text{or: } E^2 = \gamma^2 + 4\beta^2 \cos^2(k)$$

$$\text{Thus: } E = \pm \sqrt{\gamma^2 + (2\beta \cos k)^2}.$$

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go from  $[H_b, R] = 0$  to  $[H_{b'}, R^2] = 0$

- cut BZ in half
  - 2x bands
- } relabeling same states.

(half translational)

break  $R$ -trans  $\rightarrow R^2$ -trans symmetry:

- 2x bands
- $\gamma_2$  BZ
- open a gap

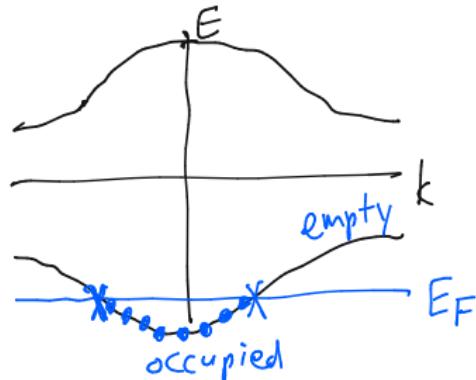
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if  $\gamma \neq 0$ :

Many mobile electrons...  
 ignore electrons...)

ground state: fill in all  
 lowest energy states

metal:  $E = E_F$  &  $E(k)$  intersect.  
 - any  $\Delta E > 0$  kicks system out of ground state.



insulator:  $E = E_F$  &  $E(k)$  not intersect

