

PHYS 4410
Quantum Mechanics 2
Spring 2023

Lecture 2

Harmonic oscillator: algebraic solution

January 23

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Last time: complicated $V(x)$

$$V(x) = \underbrace{V(x_0)}_{\text{const.}} + \underbrace{\frac{1}{2} m \omega^2 (x - x_0)^2}_{= V''(x_0)} + \dots$$

loc. of min

Convenience: - choose origin so $x_0 = 0$.
 - shift energy $V(x_0) = 0$.

energy / Hamiltonian: $H = \underbrace{\frac{p^2}{2m}}_{\text{kinetic}} + \underbrace{\frac{1}{2} m \omega^2 x^2}_{\text{potential}}$

harmonic oscillator.

Goal: e-values of H .

2 Pre-processing: "remove" as many m/ω (const.) as I can.

Idea: work in "natural units"

collect constants... determine the units.

$$[M] = \text{kg} = [M] \text{ (units of mass)}$$

$$[\omega] = \text{s}^{-1} = [T]^{-1}$$

$$[\hbar] = \text{J}\cdot\text{s} = [M][L]^2[T]^{-1}$$

Solve for $[M], [L], [T]$:

$$[M] = m$$

$$[T] = 1/\omega$$

$$[L] = \sqrt{\frac{\hbar}{m\omega}}$$

Rescale: position x :

$$\tilde{x} = \frac{x}{[L]} = \sqrt{\frac{m\omega}{\hbar}} x$$

Momentum:

$$\tilde{p} = \frac{p}{\frac{[M][L]}{[T]}} = \frac{p}{\sqrt{\hbar m \omega}}$$

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$$\begin{aligned}
 H &= \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2 \\
 &= \frac{(\sqrt{\hbar m \omega} \tilde{p})^2}{2m} + \frac{1}{2} m \omega^2 \left[\sqrt{\frac{\hbar}{m \omega}} \tilde{x} \right]^2 \\
 &= \frac{\hbar \omega}{2} [\tilde{p}^2 + \tilde{x}^2]
 \end{aligned}$$

↘ [energy]

Define $\tilde{H} = \frac{H}{\hbar \omega} = \frac{1}{2} [\tilde{p}^2 + \tilde{x}^2]$. NO "const." left.

work w/ this.

Heuristic: "work in units where $\hbar = m = \omega = 1$ "
 or... "measuring time in $1/\omega$, not s".

For today: drop ν : $\tilde{H} \rightarrow H$. Goal: $H|\psi\rangle = E|\psi\rangle$ find!

4 $H = \frac{1}{2}(p^2 + x^2) = \frac{1}{2}(x-ip)(x+ip)$ $[i = \sqrt{-1}]$

$[p = -i\hbar/dx]$ ~~if classical~~ quantum...

$$(x-ip)(x+ip) = xx - ipx + i xp + (-i^2)p^2 = x^2 + p^2 + \underbrace{i[x,p]}_{-1}$$

Claim: $[x,p] = i\hbar$

$[x,p] f(x) = x(-i\hbar \frac{d}{dx} f) - (-i\hbar \frac{d}{dx} (xf)) = i\hbar f$, so $[x,p] = i\hbar$.

Therefore: $H = \frac{1}{2}(x-ip)(x+ip) + \frac{1}{2}$

$\left\{ (x-ip)(x+ip) = 2H - 1 \right\}$

$a = \frac{x+ip}{\sqrt{2}}$ (lowering operator)

$a^\dagger = \frac{x-ip}{\sqrt{2}}$ (raising operator)

Herm. conj.

$[a, a^\dagger] = \left[\frac{x+ip}{\sqrt{2}}, \frac{x-ip}{\sqrt{2}} \right]$
 $= \frac{1}{2}(\cancel{[x,x]} + i[p,x] + \dots)$
 $= 1$

5 $H = a^\dagger a + \frac{1}{2}$ and $[a, a^\dagger] = 1$.

Suppose $H|E\rangle = E|E\rangle$.

$$\begin{aligned} H(a^\dagger|E\rangle) &= (a^\dagger a + \frac{1}{2})a^\dagger|E\rangle = (a^\dagger a a^\dagger + \frac{1}{2}a^\dagger)|E\rangle \\ &= a^\dagger(a a^\dagger + \frac{1}{2})|E\rangle \\ &= a^\dagger(a^\dagger a + [a, a^\dagger] + \frac{1}{2})|E\rangle \\ &= a^\dagger(H + 1)|E\rangle \\ &= a^\dagger(E + 1)|E\rangle = (E + 1)(a^\dagger|E\rangle) \end{aligned}$$

Also: $a|E\rangle$ is eigenstate of H w/ eigenvalue $E - 1$:

$$H(a|E\rangle) = (E - 1)(a|E\rangle).$$

6 Claim: if $H|E\rangle = E|E\rangle$, then $E \geq 0$.

$$\langle E|H|E\rangle = E \quad [\text{if } |E\rangle \text{ normalized}]$$

$$= \langle E| \frac{x^2 + p^2}{2} |E\rangle \geq 0.$$

$$\langle E|x^2|E\rangle = \int_{-\infty}^{\infty} dx \, x^2 |\psi|^2 \geq 0.$$

"large" $|E\rangle$ is in e-state w/ $E' < 0$?
 $= E - \text{"large"}$.

Resolution: must be state for which $a|\psi_0\rangle = 0$ (null vector)

$$\langle \psi_0|H|\psi_0\rangle = \underbrace{\langle \psi_0|a^\dagger a|\psi_0\rangle}_{(a|\psi_0\rangle)^\dagger (a|\psi_0\rangle) = 0} + \frac{1}{2} \underbrace{\langle \psi_0|\psi_0\rangle}_{=1} = \frac{1}{2}.$$

$$(a|\psi_0\rangle)^\dagger (a|\psi_0\rangle) = 0$$

"natural units"

So g.s. energy of H.O. $E_0 = \frac{1}{2} = \frac{1}{2} \hbar \omega$

Act w/ a^\dagger : $(a^\dagger)^n |\psi_0\rangle$ is e-state of H w/ energy $E_n = \hbar \omega (n + \frac{1}{2})$.