

**PHYS 4410**  
**Quantum Mechanics 2**  
**Spring 2023**

**Lecture 39**

**Berry phase**

May 1

**1** Adiabatic thm proof: if  $H[\lambda^\alpha(t)] \rightarrow |\psi(0)\rangle = |\psi_n(\lambda^\alpha(\omega))\rangle$   
 if  $\lambda^\alpha$ 's vary slowly enough...

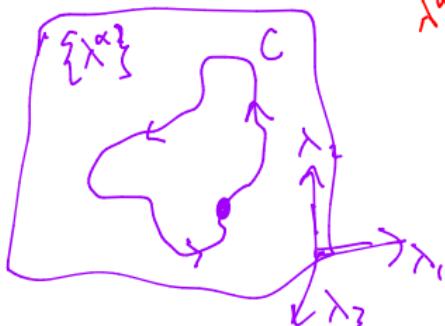
$$|\psi(t)\rangle = |\psi_n(\lambda^\alpha(t))\rangle \cdot \exp \left[ -i \int_0^t \frac{E_n(\lambda^\alpha(t'))}{\hbar} dt' \right] \cdot \exp \left[ -i \int_0^t \sum_\alpha d\lambda^\alpha \frac{dA_\alpha}{dt} \right]$$

↑  
uninteresting/time-dep PT

$$A_\alpha = -i \langle \psi_n | \frac{\partial}{\partial \lambda^\alpha} | \psi_n \rangle$$

Berry connection

Berry phase:  $\gamma = \int \sum_\alpha d\lambda^\alpha A_\alpha = \int d\vec{\lambda} \cdot \vec{A}$ , independent of time  $t$ .



$$\gamma = \oint_C d\vec{\lambda} \cdot \vec{A} \rightarrow \neq 0 \text{ in general}$$

only = 0 if  $\vec{A} = \nabla \omega$

2 Analogy with magnetism: vector potential  $\vec{A}$ :



$$\oint_C d\vec{r} \cdot \vec{A} = \oint_B = \int_R d\vec{s} \cdot \vec{B} = \text{magnetic flux}$$

$$= \int_{R'} d\vec{s} \cdot \vec{B} \quad (\text{b/c } \nabla \cdot \vec{B} = 0)$$

Take  $R$  &  $R'$  together:

$$\int_R d\vec{s} \cdot \vec{B} + \int_{R'} (-d\vec{s}) \cdot \vec{B} = \int_{\text{"sphere"}} d\vec{s} \cdot \vec{B} = \oint_C d\vec{r} \cdot \vec{A} = 0.$$

"sphere"      bdy "sphere"

gauge invariance  $\vec{B} = \nabla \times \vec{A}$ .

$$\nabla \cdot \vec{B} = \nabla \cdot (\nabla \times \vec{A}) = 0.$$

$$\vec{A} \rightarrow \vec{A} + \boxed{\nabla w(x, y, z)} : \text{leaves } \vec{B} \text{ invariant.}$$

$$\vec{A} \rightarrow \vec{A} + \nabla w :$$

$$\oint_B \rightarrow \oint_B$$

$$+ \oint_C d\vec{r} \cdot \nabla w$$

$$C$$

$$= w(\vec{r}_f) - w(\vec{r}_i)$$

3 Gauge invariance in Berry connection  $A_\alpha$ ?

$$|\psi_n(\lambda^\alpha)\rangle \rightarrow e^{i\omega(\lambda)} |\psi_n(\lambda^\alpha)\rangle$$

$$\hookrightarrow A_\alpha = -i \langle \psi_n | \frac{\partial}{\partial \lambda^\alpha} | \psi_n \rangle \rightarrow A_\alpha + \frac{\partial \omega}{\partial \lambda^\alpha}$$

Berry phase  $\gamma = \oint_C d\lambda^\alpha A_\alpha = \oint_C d\vec{\lambda} \cdot \vec{A}$  independent of  $\omega$ !

Need to generalize  $\vec{B}$  to higher d?

$$F_{\alpha\beta} = \frac{\partial A_\beta}{\partial \lambda_\alpha} - \frac{\partial A_\alpha}{\partial \lambda_\beta} = -F_{\beta\alpha}$$

↪ n-dim: Stokes Thm:



$$\sum_\alpha \oint_C d\lambda^\alpha A_\alpha = \sum_{\alpha\beta} \int_R ds^{\alpha\beta} F_{\alpha\beta} = \gamma$$

3 dimensions:

$$F = \begin{pmatrix} 0 & B_z & -B_y \\ -B_z & 0 & B_x \\ B_y & -B_x & 0 \end{pmatrix}$$

↑  
| C | ⊗ |

4 Example:  $H = \vec{\lambda} \cdot \vec{S}$

$$= \begin{pmatrix} \lambda_x & \lambda_x i \lambda_y \\ \lambda_x i \lambda_y & -\lambda_x^2 \end{pmatrix}$$

$\uparrow$   
spin- $1/2$  matrices

Parameterize:  $(\lambda_x, \lambda_y, \lambda_z) = (r \sin \theta \cos \phi, r \sin \theta \sin \phi, r \cos \theta)$

Eigenvalues of  $H$ ?

$$\det(H - E \mathbb{1}) = (E^2 - \lambda_z^2) - |\lambda_x + i \lambda_y|^2 \rightarrow E^2 - r^2 = 0.$$

$$S_x S_y + S_y S_x = 0 \quad (\text{for spin-}1/2) \quad (S_x^2 = 1) \quad \text{so } E = \pm r.$$

e.t.c.

Eigenvectors:

$$E = -r \quad H^2 = \lambda_x^2 S_x^2 \mathbb{1} + \lambda_x \lambda_y (S_x S_y + S_y S_x) + \dots = (\lambda_x^2 + \lambda_y^2 + \lambda_z^2) \mathbb{1}.$$

$$E = +r$$

$$|+\rangle = \begin{pmatrix} \cos \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} \end{pmatrix}$$

$$|-\rangle = \begin{pmatrix} -\sin \frac{\theta}{2} \\ e^{i\phi} \cos \frac{\theta}{2} \end{pmatrix}$$

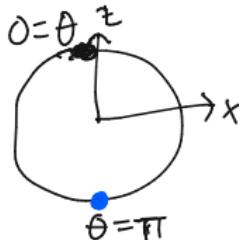
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$$E = +r$$

$$|+\rangle = \begin{pmatrix} \cos \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} \end{pmatrix}$$

$$E = -r$$

$$|-\rangle = \begin{pmatrix} -\sin \frac{\theta}{2} \\ e^{i\phi} \cos \frac{\theta}{2} \end{pmatrix}$$



Are eigenstates well-defined?

$$[0 \leq \theta \leq \pi, \quad 0 \leq \phi \leq 2\pi]$$

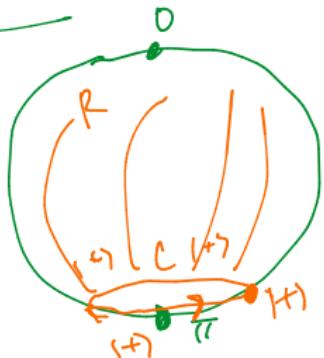
$$\theta = \pi: \quad |+\rangle = e^{i\phi} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$|-\rangle = \begin{pmatrix} -1 \\ 0 \end{pmatrix}.$$

all same physical content... but different math...

No fix!

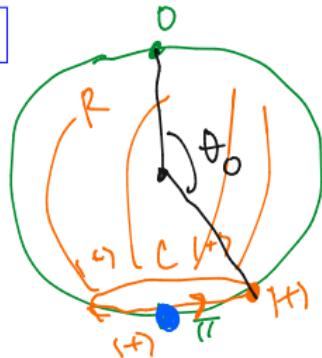
Look at  $r=1$ :



$$\gamma = \oint d\vec{x} \cdot \vec{A} = \int_R \sin \theta d\theta d\phi F_{\theta\phi}$$

$$F_{\theta\phi} = \frac{1}{2}.$$

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$$\gamma = \int_0^{\theta_0} \sin \theta d\theta \int_0^{2\pi} d\phi \cdot \frac{1}{2}$$

$$= \pi(1 - \cos \theta_0)$$

$\theta_0 \rightarrow \pi$ : loop  $C \rightarrow$  point.  $\gamma = 0$ ?

But  $\gamma = 2\pi$ .  $\xleftarrow{\text{equal!}}$

$\hookrightarrow$  what we need is  $e^{i\gamma}$ :  $e^{i0} = 1 = e^{i2\pi}$

Berry flux through any "sphere" in 3-dim parameter space:

$$\int_{\text{sphere}} F = 2\pi \cdot n$$

↑ topology/geometry  
 integer (Chern number)