

PHYS 5040
Algebra and Topology in Physics
Spring 2021

Lecture 13

February 25

1 Review the Lie group $SO(N)$, and its Lie algebra.

$$SO(N) = \text{group of rotations in } N\text{-dimensional space} \quad \frac{\text{Lie}}{\text{group}}$$
$$= \{ M \in \mathbb{R}^{N \times N} : \det(M) = 1, M^T M = \mathbb{1} \}$$

Lie algebra: $M = \mathbb{1} + \epsilon N + \mathcal{O}(\epsilon^2)$

$$\det(M) = 1 \Rightarrow \text{tr}(N) = 0 \quad M^T M = \mathbb{1} \Rightarrow \boxed{N = -N^T}$$

antisymmetric!

a basis for vector space of antisym $\mathbb{R}^{N \times N}$ matrices:

$$T_{ij} = -T_{ji} = \vec{e}_i \vec{e}_j^T - \vec{e}_j \vec{e}_i^T$$

$$i, j \in \{1, \dots, N\}$$

$$N = \sum \epsilon_{ij} T_{ij}$$

$$[T_{ij}, T_{kl}] = \delta_{jk} T_{il} - \delta_{ik} T_{jl} - \delta_{jl} T_{ik} + \delta_{il} T_{jk}$$

how many generators? $\binom{N}{2} = \frac{N(N-1)}{2}$

2 What is the transformation of a vector under rotations? What is the transformation of a tensor?

vector: $\begin{pmatrix} x_1 \\ \vdots \\ x_N \end{pmatrix} = \vec{x}$ (or $= x_i$)
 index notation $i \in \{1, \dots, N\} \in SO(N)$

N -dim representation of $SO(N) =$ fundamental

transformation of \vec{x} : $x_i \rightarrow M_{ij} x_j$

tensor: $\vec{A} \otimes \vec{B} \mapsto A_i B_j$
 rank-2 (b/c has 2 indices)
 N_{ij}

this representation $N \otimes N$

transformation of N_{ij} : $N_{ij} \rightarrow M_{ii'} M_{jj'} N_{i'j'}$

all indices are subscripts... OK for $SO(N)$
 NOT for $SU(N)$

3 What is $N \otimes N$? This gives reducible representation of $SO(N)$
 consider rank-2 tensor X_{ij} (u_i, v_j : product of 2 " N " irreps)

one irrep; $\text{tr}(X) = X_{ii}$
 transforms in the trivial representation (1)

another rep: antisymmetric rank-2 tensors: $Y_{ij} = -Y_{ji}$

Proof: $Y_{ij} + Y_{ji} = 0$

under $SO(N)$

$$M_{ii'} M_{jj'} Y_{i'j'} + M_{jj'} M_{ii'} Y_{j'i'} =$$

$$M_{ii'} M_{jj'} (Y_{i'j'} + Y_{j'i'}) = 0$$

$N=3$: irrep is vector
 $N=4$: reducible
 $N \geq 5$: irrep (new)

traceless symmetric: $X_{ij} + X_{ji} - \frac{2}{N} X_{kk} \delta_{ij} = Z_{ij}$

then $Z_{ii} = 0$ [$\delta_{ii} = N$]

irrep

4 In $SO(3)$, the antisymmetric $A_{ij} = -A_{ji}$ irrep is also **3**. Why?

$$-A_{ji} = A_{ij} \rightarrow \tilde{A}_k \quad ?$$

$SO(3)$ invariant tensors?

$$\delta_{ij} = \begin{cases} 1 & i=j \\ 0 & i \neq j \end{cases}$$

$$\begin{aligned} \delta_{ij} &\rightarrow M_{ii'} M_{jj'} \delta_{i'j'} \\ &= M_{ii'} M_{j'j} \\ &= M_{ii'} (M^T)_{i'j} \\ &= (MM^T)_{ij} = \delta_{ij} \end{aligned}$$

$$\begin{aligned} X_{ij} \delta_{ij} &\rightarrow \text{tr}(X) \\ &= X_{ii} \end{aligned}$$

is invariant under $SO(3)$

$$\boxed{\epsilon_{ijk} A_{ij} = \tilde{A}_k} \quad (j=1)$$

another invariant:
 $\epsilon_{ijk} = \begin{cases} +1 & ijk \text{ is perm. of } 123 \text{ w/ sign } = +1 \\ -1 & \text{sign } = -1 \\ 0 & 2 \text{ indices same} \end{cases}$

Why $SO(3)$ invariant?

$$\epsilon_{ijk} \rightarrow M_{ii'} M_{jj'} M_{kk'} \epsilon_{i'j'k'}$$

$$-\epsilon_{ijk} = \epsilon_{jik} \rightarrow M_{ii'} M_{jj'} M_{kk'} \underbrace{\epsilon_{j'i'k'}}_{= -\epsilon_{i'j'k'}}$$

$$\epsilon_{ijk} \xrightarrow{M} E_{ijk}$$

$$\text{but } E_{ijk} = -E_{jik} = -E_{ikj}$$

$$E(M) = \det(M) = 1 = \epsilon_{ijk} E_{ijk}$$

5 Describe all of the irreps of $SO(3)$ in terms of tensors.

$j=0$	(1) HEP	c	(scalar)	
$j=1$	(3) HEP	v_i	(vector)	$\epsilon_{ijk} X_{ij}$ $\delta_{ij} X_{ij}$
$j=2$	(5) HEP			$3 \otimes 3 = 5 \oplus 3 \oplus 1$

$$X_{ij} + X_{ji} - \frac{2}{3} \delta_{ij} X_{kk}$$

$$(j=2) \otimes (j=1) = (j=3) \oplus (j=2) \oplus (j=1)$$

$$5 \otimes 3 = 7 \oplus 5 \oplus 3$$

traceless symmetric $\leftarrow A_{ij} B_k$

$A_{ij} B_j$

$A_{ij} B_k \epsilon_{jkl}$

$+ A_{lj} B_k \epsilon_{jki}$

$A_{ij} B_k + A_{jk} B_i + A_{ki} B_j$: fully symmetric

Generalization: spin j
irrep (HEP: $2j+1$)
correspond to fully symmetric
rank j tensor

6 Describe the irreps of $SO(N)$ as tensors.

Scalar : c (1 dim)

fundamental/
vector : v_i (N dim)

traceless symmetric rank-2: $X_{ij} = X_{ji}$
 $X_{ii} = 0$ ($\frac{N(N-1)}{2} + N - 1 = \frac{(N-1)(N+2)}{2}$ dim irrep)

antisymmetric rank-2: $X_{ij} = -X_{ji}$ $\frac{N(N-1)}{2}$ dim [new irrep for $N > 4$]

Keep on going: traceless fully symmetric rank 3: $Y_{ijk} = Y_{jik} = Y_{ikj}$
 $Y_{iik} = 0$
(new irrep)

or fully antisym...
or some antisymmetry and symmetry (Young tableaux)

can use $\epsilon_{i_1 \dots i_N}$ to get rid of $\geq \frac{N}{2}$ antisymmetric indices.

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non-relativistic
The hydrogen atom (ignoring electron spin) has a "hidden" symmetry group $SO(4)$. The degeneracy of energy E_n is n^2 . Show that irreps exist of this dimension for ~~$n=1, 2, 3, 4$~~ all n . (Zee VII.11)

$|n \ell m\rangle$
 $m = -\ell, \dots, +\ell$
 $\ell = 0, \dots, n-1$
 $n = 1, 2, 3, \dots$

$$H|n \ell m\rangle = E_n |n \ell m\rangle$$

E_n expected if $SO(3)$ symmetry

Claim: fully symmetric traceless rank- n tensor, that irrep of $SO(4)$ has dimension $(n+1)^2$

$$T_{ij\dots n} = T_{ji\dots n} = T_{nj\dots i}; \quad T_{ii\dots n} = 0$$

Proof: let D_n be # of fully symmetric rank n tensors.

of traceless/sym tensor \downarrow

$$D_n = 4 + \binom{4}{2}(n-2+1) + \binom{4}{3} \frac{(n-1)(n-2)}{2} + \frac{(n-1)(n-2)(n-3)}{6}$$

$T_{11\dots 1}$ $T_{11\dots 2\dots 2}$ $T_{1\dots 2\dots 3\dots}$ $T_{1\dots 2\dots 3\dots 4\dots}$

$$C_n = D_n - D_{n-2} = (n+1)^2$$