

**PHYS 5040**  
**Algebra and Topology in Physics**  
**Spring 2021**

**Lecture 28**

April 22

**1** Review Stokes' Theorem in the language of chains and differential forms. Review closed and exact forms.

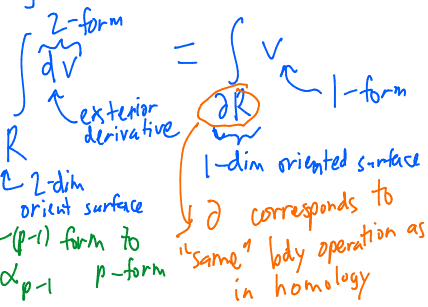
multivar calculus: Stokes' Thm:  
 $2d R \subset 3d \text{ space}$

$|U/R|$

$$\int_R (\nabla \times \vec{v}) \cdot d\vec{a} = \oint_{\partial R} \vec{v} \cdot d\vec{s}$$

in the language of differential forms:

generalize to all dimensions  
 & p-forms for any p  
 [R will be a (p+1)-dim surface]



exact differential form:  $w_p = d_{p-1} \alpha_{p-1}$   
 (p-form) (p-1 form to p-form)

closed differential form:  $d_p w_p = 0$

set of exact diff. forms:  $B_p(X)$  [X is manifold]

set of closed diff. forms:  $Z_p(X)$  so  $B_p(X) \subset Z_p(X)$

$$d_{p+1} d_p w_p = 0$$

[ $d^2 = 0$ ]

2

Write down an exact sequence involving differential forms and exterior derivatives, and define the (de Rham) cohomology groups  $H^n(X)$ .

$Z_p(X)$  &  $B_p(X)$  are Abelian groups; binary operation is addition  
 (closed) e.g. if  $w_{1,2} \in Z_p(X)$ , then  $d_p(w_1 + w_2) = d_p w_1 + d_p w_2 = 0 + 0 = 0$   
 (exact)

[identity element is 0] ;  $B_p(X) \subseteq Z_p(X)$

let's define  $F_p(X) = [\text{Abelian}]$  group of  $p$ -forms under additions  
 $[B_p, Z_p \subseteq F_p]$

$$\dots \rightarrow F_{p-1} \xrightarrow{d_{p-1}} F_p \xrightarrow{d_p} F_{p+1} \rightarrow \dots$$

$d_p$  is a homomorphism:  $d_p(w_1 + w_2) = d_p w_1 + d_p w_2$   
 $\in F_p \quad \in F_p$

And:  $d_{p+1} d_p = 0 \Rightarrow$  exact sequence [just like homology]

(de Rham) cohomology groups:

- $\text{Im}(d_p) = B_{p+1}(X)$  [ $\psi_{p+1} = d_p w_p$ ]
- $\text{Ker}(d_p) = Z_p(X)$  [ $d_p w_p = 0$ ]

$$H^p(X) := Z_p(X) / B_{p-1}(X)$$

**3** Show that if  $X$  is connected,  $H^0(X) = \mathbb{R}$ .

elements of  $H^p(X)$  are  $[w]$ , equivalence classes of "cohomologous" forms:  
 $w_p \sim w_p + d_{p-1} \lambda_{p-1}$

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Suppose:  $d_0 f = 0$ . In local coordinates  $x_i, \dots$   
scalar  $\uparrow$   
 $d_0 f = \sum_i \frac{\partial f}{\partial x_i} dx_i = 0 \Leftrightarrow \frac{\partial f}{\partial x_i} = 0 \Leftrightarrow f = \text{const.}$

$\Rightarrow$  space of cohomologous closed 0-forms = constant functions

Proof #2: let  $\gamma$  be a line segment in  $X$ .



Stokes' Thm:  $\int_{\gamma} d_0 f = \int_{\partial \gamma} f = f(b) - f(a) = 0$  if  $d_0 f = 0$

$f(b) = f(a)$  if a path connects  $a$  &  $b$ .  
constant  $\downarrow$   
 $f(x_i) = c$  if  $f$  is closed 0-form,  $H^0(X) = \{ \text{group of closed 0-forms} \} = \mathbb{R}$   
trivial no group: 1-1 form  $\downarrow$  0

4 Argue that if  $H_n(X) = \mathbb{Z}^k$ , then  $H^n(X) = \mathbb{R}^k$ .

Poincaré duality: manifold  $X$ , then:

- 1)  $H_n(X) = \mathbb{Z}^k$  for some  $k$  [no  $\times \mathbb{Z}_2$ ]
- 2)  $H^n(X) = \mathbb{R}^k$ , same  $k$

hom  $\downarrow$   
 cohom  $\uparrow$

Let  $\gamma_1, \dots, \gamma_k$  denote  $n$ -chains that are not bdy of  $(n+1)$ -chain.

$$\gamma_1 \sim \gamma_1 + \partial_{n+1} R$$

$\uparrow$   $(n+1)$ -chain

"Inner product"  $H_n \times H^n \rightarrow \mathbb{R} : \langle \gamma | \omega \rangle = \int_{\gamma} \omega$

Claim:

"inner product" is non-degenerate:  
 if  $\langle \gamma | \omega \rangle = 0 \forall \omega$ ,  
 then  $\gamma = 0$ .  
 vice versa:  
 if  $\langle \gamma | \omega \rangle = 0 \forall \gamma$ ,  
 then  $\omega = 0$ .

Note:  $\int_{\partial_p} (\omega_p + d_{p-1} \lambda_{p-1}) = \int_{\partial_p} \omega_p + \int_{\partial_p} \lambda_{p-1} = \int_{\partial_p} \omega_p$   
 $\partial_p \lambda_{p-1} = 0$

$$\int_{\gamma_p + \partial_{p+1} R_{p+1}} \omega_p = \int_{\partial_p} \omega_p + \int_{R_{p+1}} d_p \omega_p = \int_{\partial_p} \omega_p$$

$\int_{R_{p+1}} d_p \omega_p = 0$

5 For  $n = 1, 2$ , find the non-trivial element of  $H^n(S^n)$ .

$$H^0(S^1) = H^1(S^1) = \mathbb{R}$$

[NB:  $H^n(S^1) = 0$  if  $n \geq 2$  since  $d\theta \wedge d\theta = 0 \dots$ ]

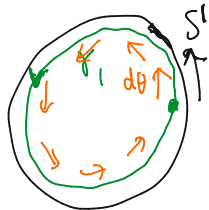
$$\rightarrow H^0(S^1) = \{f(x) = c\} \leftarrow \text{const.}$$

$H^1(S^1) = \mathbb{R}$ : look for 1-form  $\omega_1$  s.t. if  $\gamma_1$  is non-contractible loop in  $S^1$ ,

$$\int_{\gamma_1} \omega_1 \neq 0.$$

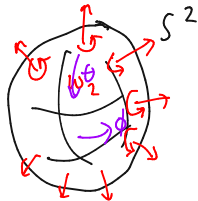
$$\int_0^{2\pi} d\theta = 2\pi \in H^1(S^1)$$

$\sim$  "length" of non-trivial closed chain  $\theta \sim \theta + 2\pi$



$$H^0(S^2) = H^2(S^2) = \mathbb{R}$$

const. ? ! ?



$\int_{S^2} \omega_2 = \text{Vol}(S^2)$  must exist. (in 2d, area...)

$$\omega_2 = \sin\theta d\theta \wedge d\phi$$

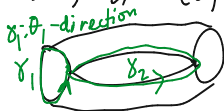
$$\int_{S^2} \omega_2 = \int_0^{2\pi} d\phi \int_0^{\pi} \sin\theta d\theta = 4\pi$$

NOT defined at  $\theta=0, \pi$

$\omega_2 \neq d(-\cos\theta d\phi)$

**6** Find all non-trivial closed but not exact forms on  $T^2$ .

"donut" or  $T^2$  :  $(\theta_1, \theta_2) \sim (\theta_1 + 2\pi, \theta_2) \sim (\theta_1, \theta_2 + 2\pi)$   
 $= S^1 \times S^1$



elements of  $H^0(T^2) = \mathbb{R}$  : 1

elements of  $H^1(T^2) = \mathbb{R}^2$  :  $c_1 d\theta_1 + c_2 d\theta_2$

elements of  $H^2(T^2) = \mathbb{R}$  :  $d\theta_1 \wedge d\theta_2$  [volume form]:

if  $\gamma_1(t) = (2\pi t, 0)$   $0 \leq t \leq 1$

$$\int_{\gamma_1} (c_1 d\theta_1 + c_2 d\theta_2) = \int_0^1 \left( c_1 \frac{d\theta_1}{dt} + c_2 \frac{d\theta_2}{dt} \right) dt = 2\pi c_1$$

$$\int_{T^2} d\theta_1 \wedge d\theta_2 = (2\pi)^2$$

- 7 Consider electromagnetism on the Riemann surface of genus  $g$ . Focusing on the vector potential  $A$ , determine the number of "pure gauge" modes that don't contribute to the magnetic field.



$$H_1(\Sigma_g) = \mathbb{Z}^{2g}$$

$$\Rightarrow H^1(\Sigma_g) = \mathbb{R}^{2g}$$

Physical DOF in magnetic field  $B = dA$   
in magnetostatics,  $B$  physical,  $A$  is not...

[not including timelike component in  $A$ ]

$$A = \underbrace{d\phi}_{\substack{\text{exact} \\ \text{1-form}}} + \underbrace{A_{\text{phys}}}_{\substack{\text{closed but not exact forms} \\ dA_{\text{phys}} \neq 0}} + A_{\text{top}}$$

from  $H^1(\Sigma_g)$

numerics:  $A_{\text{top}}$  are further redundancy than naive expectation