

**PHYS 5210**  
**Graduate Classical Mechanics**  
**Fall 2022**

**Lecture 31**

**Perturbation theory: one degree of freedom**

November 2

1 Describe qualitatively the motion of a perturbed harmonic oscillator.

$$H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2 + \underbrace{bx^4}_{\text{perturbation}}$$

small  $\downarrow$  dimensionless quantity

Solve order-by-order in small parameter ( $b$ )

$$H = \left[ \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2 \right] + \underbrace{bx^4}_{\epsilon H_1}$$

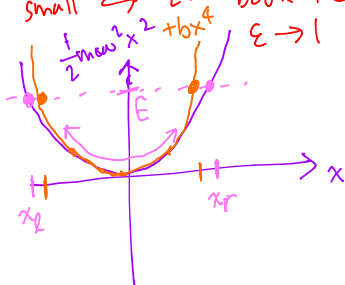
$$H = H_0 + \epsilon H_1$$

$bx^4 \ll m\omega^2 x^2$   
 Lazy all constants:  
 $b \left( \frac{E}{m\omega^2} \right)^2 \ll E$

$$b \ll \frac{(m\omega^2)^2}{E}$$

dim-less:  $\frac{bE}{(m\omega^2)^2} \} \epsilon^2$

small  $\leftrightarrow \epsilon =$  "book-keeping device"  
 $\epsilon \rightarrow 1 \quad (b \rightarrow \epsilon b)$



$(b > 0)$

expect: bounded, periodic motion

2 Solve the problem with (naive) perturbation theory.

Try:  $H = H_0 + \varepsilon H_1$

$$H_0 = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2, \quad H_1 = b x^4$$

$$x(t) = x_0(t) + \varepsilon x_1(t) + \varepsilon^2 x_2(t) + \dots$$

$$p(t) = p_0(t) + \varepsilon p_1(t) + \dots$$

Plug & chug:

$$\dot{x} = \frac{\partial H}{\partial p}$$

$$\dot{x}_0 + \varepsilon \dot{x}_1 + \dots = \left. \frac{\partial H_0}{\partial p} \right|_{p_0 + \varepsilon p_1 + \dots} + \varepsilon \left. \frac{\partial H_1}{\partial p} \right|_{p_0 + \dots}$$

$$= \frac{p}{m} = \frac{p_0 + \varepsilon p_1 + \dots}{m}$$

$$\dot{x}_0 = p_0/m$$

$$\dot{x}_1 = p_1/m \dots$$

$$\dot{x} = \frac{p}{m}$$

$$\dot{p}_0 + \varepsilon \dot{p}_1 + \dots = - \frac{\partial (H_0 + \varepsilon H_1)}{\partial x}$$

$$= -m\omega^2(x_0 + \varepsilon x_1 + \dots) - \varepsilon \cdot 4b(x_0 + \varepsilon x_1 + \dots)^3$$

$$\mathcal{O}(\varepsilon^0): \begin{cases} \dot{p}_0 = -m\omega^2 x_0 \\ \dot{x}_0 = -\omega^2 x_0 \end{cases} \left. \begin{matrix} x_0 = a \cos(\omega t) \\ \omega + \varepsilon \omega_1 \end{matrix} \right\}$$

$$\mathcal{O}(\varepsilon^1): m \ddot{x}_1 = -m\omega^2 x_1 - 4b x_0^3$$

Solution:  $x_1(t) = -\frac{A^3}{2\omega^2} \left[ \underbrace{3\omega t \sin(\omega t)}_{\cos((\omega + \varepsilon \omega_1)t)} + \frac{\cos(\omega t)}{4} - \frac{\cos(3\omega t)}{4} \right]$

blows up ☹️

**3** Why is it better to do perturbation theory with action-angle variables?  
 PT can be better behaved in other coords (not  $x, p$ )

Action-angle variables are better.

- leading order is simple.
- AA should exist for both

Caution: "hard" at higher order: see JS.

$H_0$  AND  $H_0 + \epsilon H_1$ .  
 $(J_0, \phi_0) \rightarrow (J, \phi)$   
 CT perturbatively?

Goal:  $H(J) = H_0(J, \phi) + \epsilon H_1(J, \phi)$

$$\downarrow$$

$$J = J_0 + \epsilon J_1 + \dots$$

$$\phi = \phi_0 + \epsilon \phi_1 + \dots$$

} Look for type 2 generating fcn:  
 $S(\phi_0, J)$ :

$$J_0 = \frac{\partial S}{\partial \phi_0} = J + \mathcal{O}(\epsilon)$$

$$\phi = \frac{\partial S}{\partial J} = \phi_0 + \mathcal{O}(\epsilon)$$

$$S = \underline{\phi_0 J} + \epsilon S_1 + \epsilon^2 S_2 + \dots$$

4 What is the perturbed Hamiltonian, to first order?

$$H_0(J_0) + \epsilon \underbrace{H_1(J_0, \phi_0)}_{\text{computed}} = H(J) = H(J_0 + \epsilon J_1 + \dots)$$

← dif. btwn  $J, J_0$  is  $\mathcal{O}(\epsilon)$

$$H(J) = H_0(J - \epsilon J_1 + \dots) + \epsilon H_1(J, \phi_0)$$

$$= H_0(J) - \epsilon \frac{\partial H_0}{\partial J_0} J_1 + \epsilon H_1$$

$$J_0 = \frac{\partial S}{\partial \phi_0} = J + \epsilon \frac{\partial S_1}{\partial \phi_0} + \dots = J - \epsilon J_1 + \dots, \quad \text{so } J_1 = -\frac{\partial S_1}{\partial \phi_0}$$

$$\rightarrow H(J) = H_0(J) + \epsilon \left[ H_1 + \frac{\partial H_0}{\partial J_0} \frac{\partial S_1}{\partial \phi_0} \right]$$

Goal: • only depend on  $J$ !  
• b/c  $S_1$  is unknown, use  $S_1$  to cancel  $\phi_0$  dep. in  $H_1$ !

$$H_1 = \sum_{m \in \mathbb{Z}} e^{im\phi_0} h_m(J)$$

$$S_1 = \sum_{m \in \mathbb{Z}} e^{im\phi_0} s_m(J)$$

$$\sum_{m \in \mathbb{Z}} e^{im\phi_0} \left[ h_m + \underbrace{\left( \frac{\partial H_0}{\partial J_0} \right)}_{\omega_0} im s_m \right]$$

$m=0$ : no  $\phi_0$ -dep!  
 $H = H_0 + \epsilon h_0$

$$m \neq 0: \quad s_m = -\frac{h_m}{im\omega_0}$$

5 Solve the perturbed harmonic oscillator.

$$H_0 = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2; \quad H_1 = bx^4$$

cf Lec 29?

$$H_0 = \omega J_0$$

$$x = \sqrt{\frac{2E}{m\omega^2}} \cos\phi_0 = \sqrt{\frac{2J_0}{m\omega}} \cos\phi_0$$

$$p = \sqrt{2J_0 m\omega} (-\sin\phi_0)$$

$$\begin{aligned} \text{Write } H_1(J_0, \phi_0) &= b \left( \frac{2J_0}{m\omega} \right)^2 \cos^4\phi_0 \\ &= \frac{4bJ_0^2}{m^2\omega^2} \left[ \frac{3}{8} + \frac{\cos(4\phi_0)}{8} + \frac{\cos(2\phi_0)}{2} \right] \end{aligned}$$

$$h_0 = \frac{1}{2\pi} \int_0^{2\pi} d\phi_0 H_1 = \frac{3}{2} \frac{bJ_0^2}{m^2\omega^2}$$

$$\begin{aligned} H(J) &= \omega J + \epsilon h_0(J) \quad J \approx J_0 \\ &= \omega J + \frac{3}{2} \epsilon \frac{bJ^2}{m^2\omega^2} \end{aligned}$$

What's oscillation frequency?

$$J \approx E/\omega$$

$$\omega_{\text{new}}(J) = \frac{\partial H(J)}{\partial J} = \omega + 3\epsilon \cdot \frac{bJ}{m^2\omega^2} = \omega \left[ 1 + 3\epsilon \frac{bE}{(m\omega^2)^2} \right]$$

well-behaved if  $\ll 1$