

## Homework 11

**Due:** November 17 at 11:59 PM. Submit on Canvas.

30 **Problem 1 (Adiabatic compression of a gas):** Consider a molecule in a one-dimensional non-relativistic gas. The molecule has mass  $m$  and bounces back and forth between the walls of a container of length  $L$ . The length  $L$  is assumed to slowly (adiabatically) vary in time. While in this problem we focus on the dynamics of a single molecule, we can use its behavior to estimate what happens for the gas as a whole.

1. Find an adiabatic invariant for this problem as a function of the molecule's energy,  $E$ .
2. Deduce, as  $L$  changes, how the temperature  $T$  of the gas (which is proportional to  $E$ ), will vary.
3. The pressure  $P$  in the gas is proportional to the average force per unit time applied to the left (e.g.) edge of the box, by the molecule bouncing off of it. How will  $P$  vary as a function of  $L$ ?

**Problem 2:** Consider a system with phase space  $S^1 \times \mathbb{R}$ , with canonical coordinates  $\phi$  and  $L$  where  $\{\phi, L\} = 1$ . Now consider Hamiltonian

$$H = \frac{L^2}{2} + V(\phi), \tag{1}$$

where

$$V(\phi) = \min_n |\phi - 2\pi n|. \tag{2}$$

Assume that  $\phi \sim \phi + 2\pi$  is a periodic coordinate – the form of  $V(\phi)$  is written above just to emphasize that this potential is periodic.

- 15 **A:** Sketch the  $(\phi, L)$  phase space, and sketch the trajectories on phase space generated by Hamilton's equations.
- 20 **B:** Attempt to solve this problem in action-angle coordinates. Show that you have to split the phase space into disconnected components in order to use action-angle variables, and explain why this happens (see Lecture 29). In each “patch” of phase space, find appropriate action variable  $J$  and use it to compute the period of the dynamics explicitly.

**Problem 3 (Integrability in many-body systems):** This problem will describe a curious set of differential equations that, rather remarkably: (1) turns out to be a Hamiltonian system, and (2) is integrable. The methods sketched in this problem form the starting point for a sophisticated theory of integrability in one-dimensional many-body systems.

Consider the set of differential equations

$$\dot{x}_n = x_n(x_{n+1} - x_{n-1}). \tag{3}$$

Here  $x_n$  are the degrees of freedom, and  $n$  is an arbitrary integer. You can either consider this set of equations on an infinite line, or one with periodic boundary conditions, where  $n$  is understood to be mod  $N$ . (e.g.  $x_0 = x_N$ ,  $x_1 = x_{N+1}$ , etc.).

15 **A:** This is a Hamiltonian dynamical system, if we define the Poisson bracket to be

$$\{x_n, x_{n+m}\} = x_n x_m \cdot \begin{cases} m & m = \pm 1, \\ 0 & \text{otherwise} \end{cases} \quad (4)$$

**A1.** Check that the Jacobi identity is obeyed with this Poisson bracket:

$$\{\{x_a, x_b\}, x_c\} + \{\{x_b, x_c\}, x_a\} + \{\{x_c, x_a\}, x_b\} = 0. \quad (5)$$

**A2.** Show that (3) follows from Hamiltonian mechanics, with the Hamiltonian

$$H = \sum_{n=-\infty}^{\infty} x_n, \quad \text{or} \quad \sum_{n=0}^{N-1} x_n. \quad (6)$$

20 **B:** One of the most elegant ways to understand integrability in a many-body system is to find a **Lax equation**. Let  $F$  be a *matrix* made out of the variables on phase space, and suppose that there is some other matrix  $B$  for which

$$\dot{F} = [F, B]. \quad (7)$$

**B1.** Show that (7) is obeyed if we take

$$F = \sum_n [x_n |n+1\rangle\langle n| + |n-1\rangle\langle n|], \quad (8a)$$

$$B = -\sum_n [(x_n + x_{n-1})|n\rangle\langle n| + |n-2\rangle\langle n|]. \quad (8b)$$

Here we are using bra-ket notation from quantum mechanics to denote the elements of real-valued  $N \times N$  matrices (if  $N$  is finite). There is an auxiliary basis of states labeled by the integers  $|n\rangle$ .<sup>1</sup>

**B2.** Use (7) to show that

$$Q_m = \text{tr}(F^m) \quad (9)$$

is a conserved quantity. Here you should take  $m = 2, 4, 6, \dots$  to be even.

20 **C:** Let us now show that this system is in fact integrable.<sup>2</sup> Take  $N$  finite and even.

**C1.** Explain why if  $m \leq N$ , the  $Q_m$  are linearly independent functions.

**C2.** Show that the system is integrable by explicitly checking that  $\{Q_m, Q_{m'}\} = 0$ .

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<sup>1</sup>*Hint:* Calculate  $\dot{F}|n\rangle$  and compare to  $[F, B]|n\rangle$ .

<sup>2</sup>*Hint:* I expect this part to require some mathematical creativity. My suggestion would be to start by thinking about “simple cases” and see if you tease out some general principles.