

**PHYS 5210**  
**Graduate Classical Mechanics**  
**Fall 2023**

**Lecture 10**  
**Euler's equations**

September 20

$$RR^T = I$$

orthogonal  
 $[3 \times 3^\vee$  matrices  
 $\det = 1]$

Configuration space of rigid body rotation:  $SO(3)$  [

Physically: matrix  $R$  is rotation from **body frame**  $\xrightarrow{\hspace{2cm}}$  **space frame**  
 $(I) \qquad \qquad \qquad (i)$

$$R_{iI} : \left( \sum_{I=1}^3 R_{iI} R_{jI} \right) = \delta_{ij} = \begin{cases} 1 & i=j \\ 0 & i \neq j \end{cases}. \quad R_{iI} R_{iJ} = \delta_{IJ}$$

Lagrangian for rigid body dynamics:

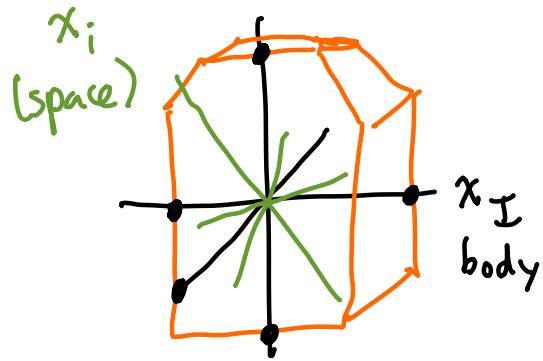
$$L = f(R_{iI}, \dot{R}_{iI}, \dots) + \Lambda_{IJ} (R_{iI} R_{iJ} - \delta_{IJ})$$

Lagrange multiplier:  $\Lambda_{IJ} = \Lambda_{JI}$

Use effective theory to constrain  $f$ :



## ① Symmetries:



A rigid body may appear asymmetric in body frame

No preferred orientation in space.  
(Choice of "origin", or "identity matrix", is arbitrary)

Physical: relative rotation over time.

Intuitively: like translation? (no preferred  $x=0$ ), but on  $SO(3)$ ?  
orthogonal/ $SO(3)$

Define: left  $SO(3)$  invariance:  $S[R_{iI}(t)] = S[Q_{ij} \downarrow R_j I^{(+)})]$   
for any constant  $Q_{ij}$ . [Generalizes:  $x(t) \rightarrow x(t) + \varepsilon$ ]

Contrast: w/ right  $SO(3)$ :  $S[R_{iI}(t)V_{IJ}] \stackrel{?}{=} S[R_{iJ}(t)]$   
NOT symmetry; unless object is sphere.

Symmetries of rigid body rotation:

- left  $SO(3)$  invariance
- time-reversal symmetry
- right symmetry: of object

## ② Invariant BBS:

Under time-reversal:  $t \rightarrow -t$

$$\dot{R}_{iI} \rightarrow -\dot{R}_{iI}$$

↓  
2 time derivatives in  $L$

Under left  $SO(3)$  invariance:  
invariants of  $SO(3)$ :  $\delta_{ij}$ ,  $\epsilon_{ijk}$  not use

Analogy to relativity:  $\eta_{\mu\nu}$  invariant under Lorentz.

$\eta_{\mu\nu} a^\mu b^\nu$  invariant

→ contract indices w/  $\delta_{ij}$ : (no dangling space frame indices)

$$\cancel{R_i I R_{iJ}}$$

$\delta_{IJ} = \text{const.}$   
not dynamical

$$\cancel{R_i I R_{iJ}}$$

$\cancel{\Omega_{IJ}}$

$$\underbrace{R_i I R_{iJ}}_{\text{nice invariant.}}$$

Under right symmetry (body frame symmetry):

Today: no symmetry in body frame;

$$L = \frac{1}{2} \dot{R}_{iI} \dot{R}_{iJ} K_{IJ} + \Lambda_{IJ} (R_{iI} R_{iJ} - \delta_{IJ})$$

Symmetric matrix

has manifest left- $SO(3)$  invariance.

③ EOMs:

$$\frac{\delta S}{\delta \dot{R}_{IJ}} = 0, \quad \text{or} \quad R_{iI} R_{iJ} = \delta_{IJ}$$

$$\frac{d}{dt} (R_{iI} R_{iJ}) = 0$$

$$= \dot{R}_{iI} R_{iJ} + R_{iI} \dot{R}_{iJ}$$

$$= R_{iJ} \dot{R}_{iI} + R_{iI} \dot{R}_{iJ} = \underbrace{\Omega_{JI} + \Omega_{IJ}}_{\Omega_{JI} + \Omega_{IJ} = 0} = 0$$

antisymmetric  
 $\Omega = -\Omega^T$

$\dot{R}_{ij} = R_{ij} \Omega_{ij}$ ;  $\Omega_{ij}$  represents angular velocities relative to body frame!

$$\frac{\delta S}{\delta R_{ij}} = 0 = -\frac{d}{dt}(\dot{R}_{ij} K_{ij}) + \Lambda_{ij} R_{ij}$$

$$\ddot{R}_{ij} K_{ij} R_{ik} = \Lambda_{ik}$$

$$\ddot{R}_{ij} = \frac{d}{dt}(R_{il} \Omega_{lj}) = \dot{R}_{il} \Omega_{lj} + R_{il} \dot{\Omega}_{lj}$$

$$\text{so: } \dot{R}_{ij} K_{ij} R_{ik} = \underline{\Omega_{kl} \Omega_{lj} K_{ij}} + \underline{\dot{\Omega}_{kj} K_{ij}} = \Lambda_{ik}$$

Get rid of  $\Lambda_{ik} = \Lambda_{ki}$  by using symmetric: take antisymmetric part!

$$0 = \dot{\Omega}_{kj} K_{ij} - \dot{\Omega}_{ij} K_{kj} + \Omega_{kl} \Omega_{lj} K_{ij} - \Omega_{il} \Omega_{lj} K_{kj}$$

Reduced dynamical system to body frame only. [left  $SO(3)$  invariance]

Define angular momenta:

$$L_{ij} = -L_{ji} = K_{ik} \Omega_{kj} + \Omega_{ik} K_{kj}$$

$$[L = K\Omega + \Omega K]$$

↑  
sym      antisym

$$\text{Show: } \dot{L}_{ij} = L_{ik} \Omega_{kj} - \Omega_{ik} L_{kj} = [L, \Omega]_{ij}$$

Write out:

angular velocity around 3-axis of body frame

$$L_{ij} \rightarrow \begin{pmatrix} 0 & -L_3 & L_2 \\ L_3 & 0 & -L_1 \\ -L_2 & L_1 & 0 \end{pmatrix}$$

$$\Omega_{ij} \rightarrow \begin{pmatrix} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{pmatrix}$$

$\dot{L} = [L, \Omega] \rightarrow$  Euler's equations:

$$\boxed{\begin{aligned}\dot{L}_1 &= L_2 \omega_3 - L_3 \omega_2 \\ \dot{L}_2 &= L_3 \omega_1 - L_1 \omega_3 \\ \dot{L}_3 &= L_1 \omega_2 - L_2 \omega_1\end{aligned}}$$

Convenient to choose basis in body frame where  
 $K_{IJ}$  is diagonal.

Define (principal) moments of inertia:

$$\begin{aligned}I_1 &= K_2 + K_3 \\ I_2 &= K_3 + K_1 \\ I_3 &= K_1 + K_2\end{aligned}$$

$$K = \begin{pmatrix} K_1 & 0 & 0 \\ 0 & K_2 & 0 \\ 0 & 0 & K_3 \end{pmatrix}$$

$$\left. \begin{aligned}L_1 &= I_1 \omega_1 \\ L_2 &= I_2 \omega_2 \\ L_3 &= I_3 \omega_3\end{aligned} \right\} L = K \Omega + \Omega K$$

$$\hookrightarrow \text{Euler: } \begin{aligned}I_1 \dot{\omega}_1 &= (I_2 - I_3) \omega_2 \omega_3 \\ I_2 \dot{\omega}_2 &= (I_3 - I_1) \omega_3 \omega_1 \\ I_3 \dot{\omega}_3 &= (I_1 - I_2) \omega_1 \omega_2.\end{aligned}$$