

**PHYS 5210**  
**Graduate Classical Mechanics**  
**Fall 2023**

**Lecture 19**

**Elastic solids**

October 11

Degrees of freedom in solid:

$\sigma^I(x_i, t)$  = equilibrium location of piece of solid is now at  $x_i$  at time  $t$ .  $I=1,2,3$ .

$$\mathcal{L} = \frac{\rho}{2} \partial_t \sigma^I \partial_t \sigma^I - \frac{1}{8} \lambda^{IJKL} (\partial_i \sigma^I \partial_j \sigma^J - \delta^{IJ}) (\partial_j \sigma^K \partial_l \sigma^L - \delta^{KL})$$

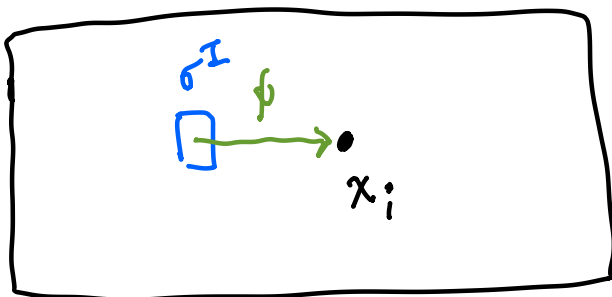
$\lambda^{IJKL} = \tilde{K} \delta^{IJ} \delta^{KL} + \mu (\delta^{IK} \delta^{JL} + \delta^{IL} \delta^{JK})$

for isotropic solid,  $\lambda$  can only contain  $\delta$ 's  $\rightarrow$  HW7 relaxes.

Today: behavior of solids close to equilibrium.

$$\sigma^I = \delta^I_i (x_i - \phi_i)$$

$\phi$  infinitesimally small,  
only keep  $\mathcal{O}(\phi)$  terms  
in EOMs.



Expand  $\mathcal{L}$  to quadratic order in  $\phi$ :

$$\bullet \partial_t \sigma^I = \delta^I_i (-\partial_t \phi_i)$$

$$\begin{aligned} \bullet \partial_i \sigma^I \partial_j \sigma^J - \delta^{IJ} &= \partial_i \left[ \delta^I_m (\pi_m - \phi_m) \right] \partial_j \left[ \delta^J_n (\pi_n - \phi_n) \right] - \delta^{IJ} \\ &= \delta^I_m \delta^J_n (\delta_{im} \delta_{jn} - \delta_{im} \partial_i \phi_n - \delta_{jn} \partial_j \phi_m + \cancel{\partial_i \phi_m \partial_j \phi_n}) - \delta^{IJ} \\ &= \cancel{\delta^{IJ}} - \cancel{\delta^{IJ}} - \delta^I_m \delta^J_n (\partial_m \phi_n + \partial_n \phi_m) + \text{quadratic } \phi \end{aligned}$$

↓  $\frac{\partial_m \phi_n + \partial_n \phi_m}{2}$

Strain tensor:  $u_{mn} = \frac{\partial_m \phi_n + \partial_n \phi_m}{2}$

$$\begin{aligned} \mathcal{L} &\approx \frac{\rho}{2} \partial_t \phi_i \partial_t \phi_i - \frac{1}{8} \lambda^{IJKL} \left[ \delta^I_n \delta^J_n (-\partial_m \phi_n - \partial_n \phi_m) \right] \left[ \delta^K_p \delta^L_q (-\partial_p \phi_q - \partial_q \phi_p) \right] \\ &= \frac{\rho}{2} \partial_t \phi_i \partial_t \phi_i - \frac{1}{8} \lambda_{mnpq} (\partial_m \phi_n + \partial_n \phi_m) (\partial_q \phi_p + \partial_p \phi_q) \end{aligned}$$

(Vestigial) symmetries?

-  $x_i$  &  $t$  translation  $[x_i \rightarrow -x_i \quad t \rightarrow -t \text{ also!}]$

-  $\phi_i \rightarrow \phi_i + \varepsilon_i$  shift symmetry [combo of  $\pi$  &  $\sigma$  shift]

global displacement of solid OK.

-  $x_i$  rotation symmetry:

① if  $\lambda_{mnpq} = \tilde{K} \delta_{mn} \delta_{pq} + \dots$  then  $x_i \rightarrow Q_{ij} x_j$

②  $\mathcal{L} = 0$  if  $\phi_i = \underline{\varepsilon_{ij}} x_j$   $\varepsilon_{ij} = -\varepsilon_{ji}$   
antisym.

$$\partial_i \phi_j + \partial_j \phi_i = \varepsilon_{ji} + \varepsilon_{ij} = 0.$$

↳  $-\sigma^I$  rotation symmetry  $\rightarrow$  symmetries of  $\lambda_{ijkl}$

-  $\phi_i \rightarrow \varepsilon_{ij} x_j$  encodes space frame

$x_i \rightarrow Q_{ij} x_j$  rotation of  $\mathcal{L}(\sigma^I)$ .

Physical predictions of theory:

- normal modes (lec. 20)
- **stress** on solids (time-independent)

Calculate stress tensor from Noether's Thm.

lec 15: stress tensor from  $x_i$  translation... in  $\mathcal{L}(\sigma, \kappa)$   
[not  $\mathcal{L}(\phi, \kappa)$ ].

$$T_{ij} = \mathcal{L} \delta_{ij} - \partial_j \sigma^I \frac{\partial \mathcal{L}}{\partial (\partial_i \sigma^I)}$$

$$= \mathcal{L} \delta_{ij} - \underbrace{\partial_j \sigma^I}_{\approx \delta_j^I(\phi^0)} \left[ -\frac{1}{4} (\lambda^{IJKL} + \lambda^{JIKL}) \partial_i \sigma^J \underbrace{(\partial_j \sigma^K \partial_j \sigma^L - \delta^{KL})}_{\phi + \phi^2} \right]$$

Lowest order in  $\phi$ ? Linear in  $\phi$ ...

$$T_{ij} \approx \delta_j^I \frac{1}{2} \lambda^{IJKL} \delta_i^J (-\delta_m^K \delta_n^L) (\partial_m \phi_n + \partial_n \phi_m)$$

$$= -\frac{1}{2} \lambda_{ijkl} (\partial_k \phi_l + \partial_l \phi_k) = -\lambda_{ijkl} u_{kl}$$

Theory of elastic solids:

stress (T) linearly prop to strain (u)

When is equilibrium stable? " $\lambda_{ijkl} \geq 0$ "

[analogue to Hooke's Law:  $F = -kx$ ,  $k > 0$ ].

Isotropic solid:  $\lambda_{ijkl} = \tilde{K} \delta_{ij} \delta_{kl} + \mu (\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk})$

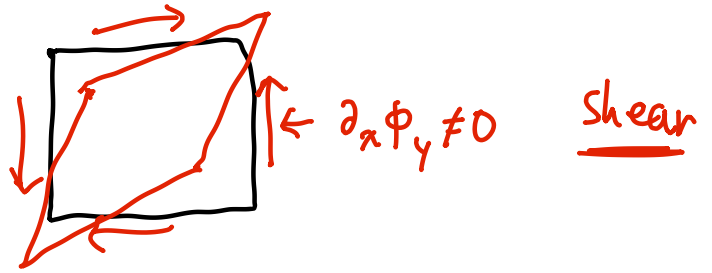
← Lamé constants / elastic moduli;

Potential energy:  $\sim \frac{1}{2} \lambda_{ijkl} u_{ij} u_{kl} \geq 0$  [term in  $\mathcal{L}$ ]

Check this positivity:

①  $u_{xy} = u_{yx} \neq 0$

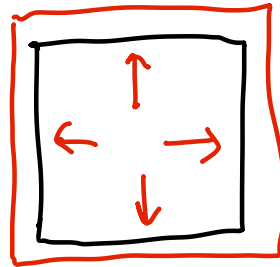
$$u = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$



$$\lambda_{ijkl} u_{ij} u_{kl} \geq 0$$

$$= \lambda_{xyxy} + \lambda_{xyyx} + \lambda_{yxxy} + \lambda_{yxyx} = 4\mu \rightarrow \mu \geq 0$$

②  $u_{ij} = \delta_{ij}$



expansion/  
compression

$$\lambda_{ijkl} \delta_{ij} \delta_{kl} \geq 0$$

$$\hookrightarrow \underbrace{\tilde{K} \delta_{ij} \delta_{ij} \delta_{kl} \delta_{kl}}_{\text{tr}(D)=3} + \mu \delta_{ij} \delta_{kl} (\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk}) = 9\tilde{K} + 6\mu \geq 0$$

or  $\tilde{K} \geq -\frac{2}{3}\mu$ .

Write  $K = \tilde{K} + \frac{2}{3}\mu$  so  $K \geq 0$ .

$$\lambda_{ijkl} = \underbrace{K}_{\text{bulk modulus}} \delta_{ij} \delta_{kl} + \underbrace{\mu}_{\text{shear modulus}} \left( \delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk} - \frac{2}{3} \delta_{ij} \delta_{kl} \right)$$

$\delta_{ij} \delta_{kl}$  = projection onto spin 0 of trace

$\hookrightarrow$  projection onto spin 2 shearing solid / volume-preserving deform. vector

Rotation invariance:  $3 \times 3$  matrix = (spin 1)  $\oplus$  (spin 1)

$=$  (spin 0)  $\oplus$  (spin 1)  $\oplus$  (spin 2)

$\lambda_{ijkl} u_{kl} = K \delta_{ij} \text{tr}(u)$  (trace)

traceless symmetric antisymmetric  $\rightarrow$  global rotation of solid.