

PHYS 5210
Graduate Classical Mechanics
Fall 2023

Lecture 28
Hamilton-Jacobi equation

November 1

lec 27: Type I CT: $(x_i, p_i) \rightarrow (X_i, P_i)$ where
 (x_i, X_i) form independent coords on phase space.

Generating function F_1 ; $dF_1 = p_i dx_i - P_i dX_i$ ($p_i = \frac{\partial F_1}{\partial x_i}$)

Hamiltonian for X, P, \dots : $H' = H + \frac{\partial F_1}{\partial t}$

Can we choose F_1 so that:

① $H' = 0$.

② (X_i, P_i) to be constants of motion.

① \Rightarrow ②: $\dot{X}_i = \frac{\partial H'}{\partial p_i} = 0$ $\dot{P}_i = -\frac{\partial H'}{\partial X_i} = 0$.

Each $(X_i, P_i) \rightarrow$ one choice of init cond for Ham eq
in x_i, p_i .

If so... find $x_i(X, P, t)$ and $p_i(X, P, t)$.

Such a special F_1 will be called S :

Then: $0 = H' = \frac{\partial S}{\partial t} + H(x_i, p_i, t)$
 $\hookrightarrow = \frac{\partial S}{\partial x_i}$

Hamilton-Jacobi: $0 = \frac{\partial S}{\partial t} + H(x_i, \frac{\partial S}{\partial x_i}, t) \rightarrow n$ coords x_i, t

If solved...

$S(x_i, X_i, t)$ depends on integration constants X_i

Problem: HJ is a partial differential equation
... much harder than ODE! $\hat{=}$

Strategy: look for special solutions to HJ, enabled by
"separation of variables"

Example 1: central force problem:

$$L = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2) - V(r)$$

Legendre transform!

$$H = \frac{p_r^2}{2m} + \frac{p_\theta^2}{2mr^2} + V(r).$$

go to HJ

$$0 = \frac{\partial S}{\partial t} + \frac{1}{2m} \left(\frac{\partial S}{\partial r} \right)^2 + \frac{1}{2mr^2} \left(\frac{\partial S}{\partial \theta} \right)^2 + V(r)$$

"Separation of variables":

$$S(r, \theta, t) = S_r(r) + S_\theta(\theta) + S_t(t)$$

$$0 = \dots + \frac{1}{2mr^2} S_\theta'(\theta)^2 \rightarrow \text{bad! both } r \text{ \& } \theta \text{ dep.} \dots$$

{ Aside: s.o.v. for Laplace: $\phi(x, y) = X(x)Y(y)$:

$$\frac{1}{\phi} \nabla^2 \phi = \frac{X''(x)}{X} + \frac{Y''(y)}{Y} = 0. \quad \frac{X''}{X} = \lambda \quad \frac{Y''}{Y} = -\lambda.$$

Need $\frac{\partial S}{\partial t} = \text{const.}$ since no other t -dep.

So: $S_t = -Et$, AND $S_\theta = L\theta$

↳ Plug into HS; $0 = -E + \frac{S_r'^2}{2m} + \frac{L^2}{2mr^2} + V(r)$

consistent since each term only depends on r .

$$S_r'^2 = 2mE - \frac{L^2}{r^2} - V(r) \cdot 2m$$

Solve up to quadratures (HW 10)

or $S_r(r) = \int dr \sqrt{2mE - \frac{L^2}{r^2} - 2mV(r)}$

We've found: $S(r, \theta, t; E, L)$ "integration constants"

Idea: treat E, L as our X_i .

Now find other 2 const. of motion:

$$-P_E = \frac{\partial S}{\partial E} = -t + \frac{\partial S_r}{\partial E} = -t + \int \frac{dr \cdot 2m}{2 \sqrt{2mE - L^2/r^2 - 2mV}}$$

$$-P_L = \frac{\partial S}{\partial L} = \theta + \frac{\partial S_r}{\partial L} = \theta - \int \frac{dr \cdot L/r^2}{\sqrt{2mE - 2mV - L^2/r^2}}$$

Example 2: particle in magnetic field

$$H = \frac{p_x^2}{2m} + \frac{(p_y + qBx)^2}{2m}$$

$$\hookrightarrow 0 = \frac{\partial S}{\partial t} + \frac{1}{2m} \left(\frac{\partial S}{\partial x} \right)^2 + \frac{1}{2m} \left(\frac{\partial S}{\partial y} + qBx \right)^2$$

Separate variables: $S = -Et + p_y y + W(x)$:

$$2mE = W'(x)^2 + (p_y + qBx)^2$$

so: $W(x) = \int dx \sqrt{2mE - (p_y + qBx)^2}$.

Find new conserved quantities:

$$\frac{\partial S}{\partial E} = -t + \int dx \frac{2m}{2\sqrt{2mE - (p_y + qBx)^2}}$$

$$= -t - m \int \frac{\sqrt{2mE} \sin\theta d\theta}{qB \sqrt{2mE} (1 - \cos^2\theta)}$$

$$\text{const.} = -t - \frac{m}{qB} \theta$$

$$\theta = -\frac{qB}{m}(t - t_0) \rightarrow \cos\left(\frac{qB}{m}(t - t_0)\right) = \frac{p_y + qBx(t)}{\sqrt{2mE}}$$

Then:

$$\text{const.} = \frac{\partial S}{\partial p_y} = y - \int dx \frac{p_y + qBx}{\sqrt{2mE - (p_y + qBx)^2}}$$

$$= y + \frac{\sqrt{2mE}}{qB} \sin\theta$$

or...

$$\underbrace{(x - x_0)^2 + (y - y_0)^2}_{\text{const.}} = \frac{2mE}{(qB)^2} \rightarrow \text{circular trajectories}$$