

PHYS 5210
Graduate Classical Mechanics
Fall 2023

Lecture 31

Integrable systems

November 8

Hamiltonian H is integrable if there exist n independent conserved I_A ($A=1, \dots, n$) on $2n$ -dim phase space. Obey:

(in involution) $\{I_A, I_B\} = 0$ for all A and B .

Always take $I_1 = H$ (if H is t -ind.)

Liouville's Integrability Thm: if surfaces of const. I_A are compact (don't go to ∞), then there are action-angle variables

(ϕ_A, J_A) w/ $\{\phi_A, J_B\} = \delta_{AB}$ and $H = H(J)$.

A-A vars completely solve problem:

$$\dot{J}_A = -\frac{\partial H}{\partial \phi_A} = 0$$

$$J_A = \text{const.}$$

$$\text{and } \dot{\phi}_A = \frac{\partial H}{\partial J_A} = \omega_A = \text{const.}$$

$$\phi_A(t) = \phi_A(0) + \omega_A t$$

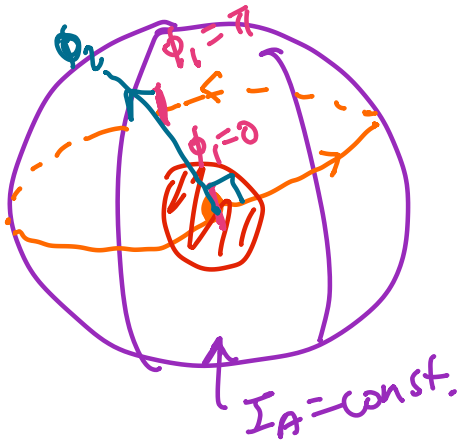
Thus integrable system is:

① exactly solved

② gives geometric picture for phase space.

Sketch LI Theorem's proof:

① By lec 24, each I_A generate CTs along surface of const. I_A .
 (b/c $\delta I_A = \epsilon \{I_A, I_A\} = 0$)



Claim: there exists a $J_1(I_1, \dots, I_n)$ where flow returns to starting point. (in a local region)

Flow needs to return near start b/c surface is compact.

Then define angle variable ϕ_1 to be coord. along flow

Normalize ϕ_1 & J_1 so that $\phi_1 \sim \phi_1 + 2\pi$.

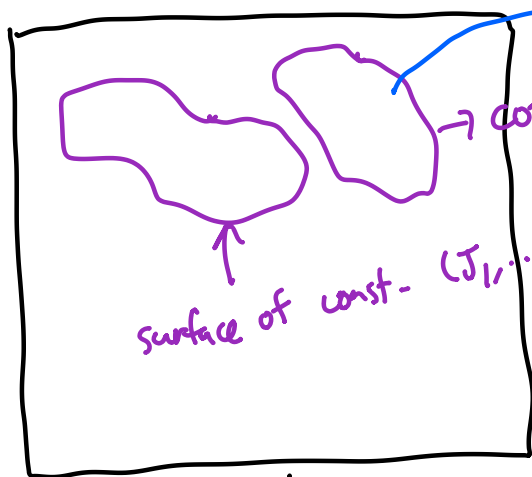
② Iterate... Choose $J_2(I_1, \dots, I_n)$ indep. of J_1 and choose to generate flow perp. to ϕ_1 ... give us ϕ_2

These flows preserve all I_A & J_A ... and

$$\{\phi_1, J_2\} = 0, \quad \{\phi_2, J_1\} = 0.$$

Normalize so $\{\phi_2, J_2\} = 1$, etc....

We can do this construction if I_A 's are smooth...



phase space

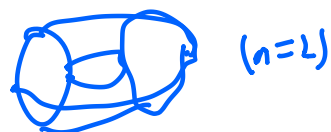
LI Thm:



space "spanned" by ϕ_A

= n-dim. torus (T^n)

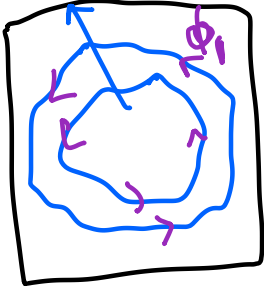
= $S^1 \times S^1 \times \dots \times S^1$



(n=2)

"artist rendition"

If $n=1$: J_1



A-A as useful "polar" coords...

LI Thm: keeps working in $2n$ -dim phase space.

This is abstract... in practice:

- find AA vars by separating HJ eq?
- or stop w/ I_A 's ... (see HW 11)

Well developed theory of integrability. --

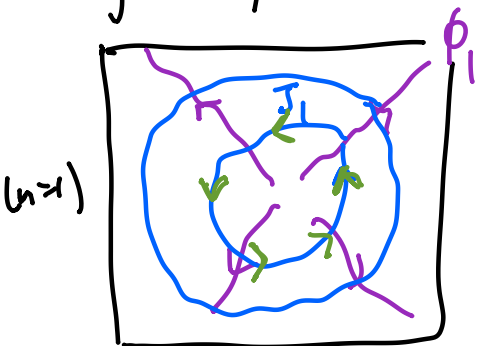
strongest in one spatial dim (lattice / field)
not much for $d > 1$ except harmonic osc.

How can something NOT be integrable?

"Loophole" = I_A 's might not be smooth (lec 34+)

Much of our knowledge of physics based on integrable systems, or perturbations of them.

Integrability:



$$\dot{J}_A = 0$$

$$\dot{\phi}_A = \frac{\partial H}{\partial J_A} = \omega_A = \text{const.}$$

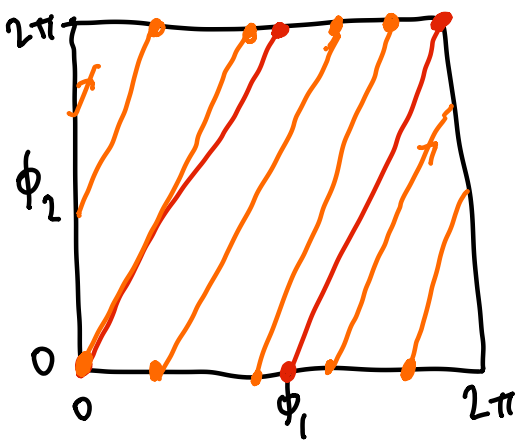
wind around each circle of T^1
at frequency ω_A

Now: $n=2$?

Draw T^2 as box w/ opposite edges ident.

① commensurate: $\frac{\omega_1}{\omega_2} = \text{rational} (= \frac{1}{2})$
(ratio of integers)

② incommensurate: $\frac{\omega_1}{\omega_2} = \text{irrational}$



Commensurate: trajectory = one-dimensional closed loop.

Incom: never return to origin, trajectory "fills up T^2 ".

for each point (ϕ_1, ϕ_2) , we get w/in dist ϵ of point for any $\epsilon > 0$. Trajectory is "dense"

math fact:

$p + q \cdot \alpha$ can be as close as we want to any real $\neq \pi$.
↑ ↑
integers irrational (fix $\alpha!$)

can be as close as we want to any real $\neq \pi$.

"ergodicity"

$$= \frac{1}{T} \int_0^T dt F(t) \rightarrow \text{avg of } F \text{ on } T^n.$$
$$= \frac{1}{(2\pi)^2} \int d\phi_1 d\phi_2 F$$

Incom: seems like 1d trajectory \rightarrow exploring 2 dim space?

Integrable system can traverse 1, 2, 3, ..., n-dim subspace of phase space.

Chaotic systems can explore 1.76...-dimensional subspace (fractals!)