

PHYS 5210
Graduate Classical Mechanics
Fall 2023

Lecture 35

Kicked rotor

November 17

lec 34: t -ind $H \rightarrow 2$ DOF for chaos

Today: t -dep $H \rightarrow 1$ DOF for chaos

Example: kicked rotor. let $\{\phi, J\} = 1$.

$$\phi \sim \phi + 2\pi$$

periodically identify

$$H(t) = \frac{J^2}{2I} + \sum_{n=-\infty}^{\infty} \delta(t - n\tau) \cdot \varepsilon_0 \cos \phi$$

Dirac δ function.

EOMs:

$$\dot{\phi} = \frac{\partial H}{\partial J} = \frac{J}{I}$$

$$\dot{J} = -\frac{\partial H}{\partial \phi}$$

$$= \sum_{n=-\infty}^{\infty} \delta(t - n\tau) \varepsilon_0 \sin \phi$$

$\dot{J} = 0$ if $t \neq n\tau$.

so $J(t)$ will be piecewise const.

Let's suppose at $t=0^+$ (just after $t=0$):

$$\phi(0^+) = \phi_0, \quad J(0^+) = J_1 \quad \leftarrow \text{explain later!}$$

Since $\dot{J}=0$ for $0 < t < \tau \longrightarrow J(t) = J_1 \quad (0 < t < \tau)$.

And $\dot{\phi} = \frac{J}{I}$ means $\phi(t) = \phi_0 + \frac{J_1}{I} t \quad (0 < t < \tau)$.

Just before $t=\tau$ kick: $\phi(\tau^-) = \phi_1 = \phi_0 + \frac{\tau}{I} J_1 \quad \textcircled{1}$

What about after kick?

$$J(\tau+\delta) = J(\tau-\delta) + \int_{\tau-\delta}^{\tau+\delta} dt \dot{J}(t) = J(\tau-\delta) + \int_{\tau-\delta}^{\tau+\delta} dt \sum_{n=-\infty}^{\infty} \delta(t-n\tau) \epsilon_0 \sin \phi(t)$$

$$= J_1 + \epsilon_0 \sin \phi(\tau) \quad \textcircled{2}$$

$\int_{\tau-\delta}^{\tau+\delta} dt \delta(t-\tau) f(t) = f(\tau)$
 [other $n \neq 1$ terms $\rightarrow 0$]

$$J(\tau^+) = J_2 = J_1 + \epsilon_0 \sin \phi_1$$

$$\leq \delta \cdot \frac{1}{I} \max(J_1, J_2) \rightarrow 0.$$

Justify $\phi(\tau) = \phi_1$?

$$\phi(\tau+\delta) = \phi(\tau-\delta) + \int_{\tau-\delta}^{\tau+\delta} dt \dot{\phi} = \phi(\tau-\delta) + \int_{\tau-\delta}^{\tau+\delta} dt \frac{J(t)}{I}, \quad \text{so } \phi(t) \text{ is continuous}$$

Iterate this argument, define

$$\phi(n\tau^-) = \phi_n, \quad J(n\tau^-) = J_n.$$

$$\textcircled{2} \rightarrow J_{n+1} = J_n + \epsilon_0 \sin \phi_n$$

$$\textcircled{1} \rightarrow \phi_{n+1} = \phi_n + \frac{\tau}{I} J_{n+1}$$

} reduces Hamilton's eqs (ODEs) to discrete map.

Reduced to math... work in dimensionless units!

$$[\phi] = \text{const.}$$

$$[t] = [T]$$

$$[J] = [M][L]^2[T]^{-1}$$

mass length

$$[\tau] = [T]$$

$$[\varepsilon_0] = [J] = [M][L]^2[T]^{-1}$$

$$[I] = [M][L]^2$$

Define $\tilde{\phi}_n = \phi_n$ (dim-less), $\tilde{J}_n = \frac{J_n}{I/\tau}$

$$\varepsilon = \frac{\varepsilon_0 \tau}{I}$$

dimensionless kick strength

Dimensionless kicked rotor:

$$\tilde{J}_{n+1} = \tilde{J}_n + \varepsilon \sin \tilde{\phi}_n$$

$$\tilde{\phi}_{n+1} = \tilde{\phi}_n + \tilde{J}_{n+1}$$

NOTE: $\tilde{\phi}_n + 2\pi \rightarrow \tilde{\phi}_n$ (same pt in phase space)

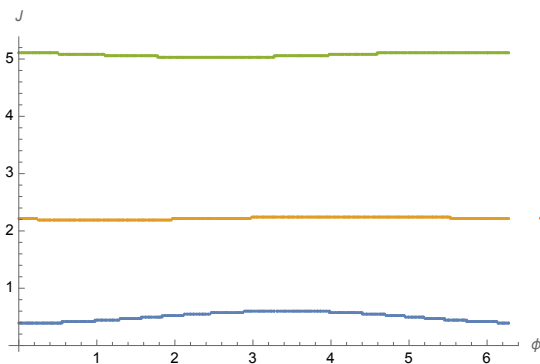
$$\tilde{J}_n + 2\pi \rightarrow \tilde{J}_n$$

doesn't change any $\tilde{\phi}_n \dots$

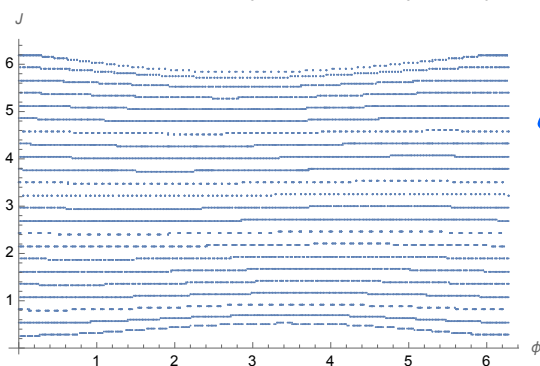
going to identify $\tilde{J} \sim \tilde{J} + 2\pi$.

Henceforth, drop tildes...

Numerical simulations: $\varepsilon = 0.05$



← plot of 3 diff. ICs (each colour)
parametric plot (ϕ_n, J_n)

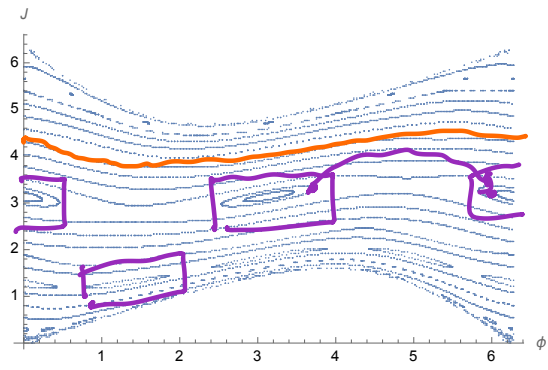
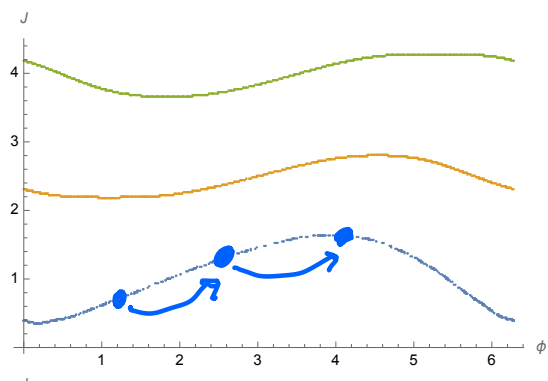


← 20 diff ICs.

Punchline: each trajectory seems confined to 1d?

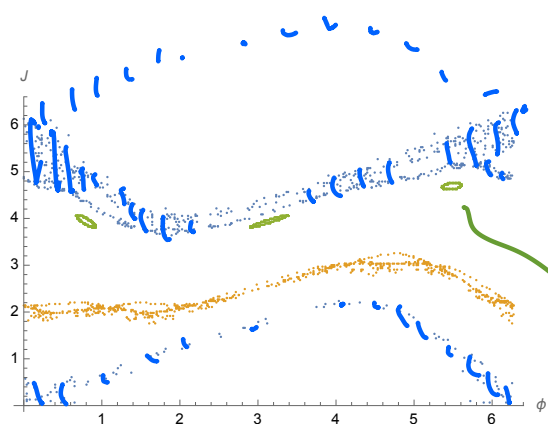
Dynamics looks integrable: motion restricted to 1d?

Numerics at $\epsilon = 0.6$.
 still integrable?



Some "trajectories" (ϕ_n, J_n)
 wind around ϕ .

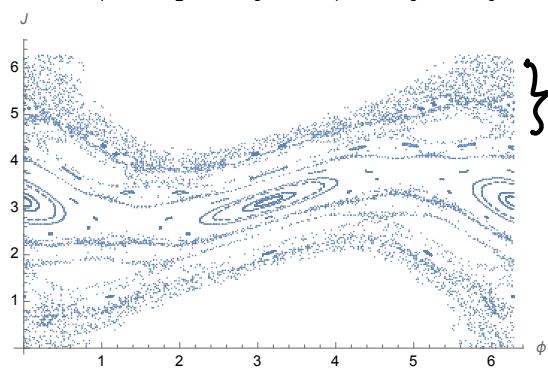
trapped "orbits" also exist.



Numerics at $\epsilon = 1.05$.

→ chaotic dynamics?

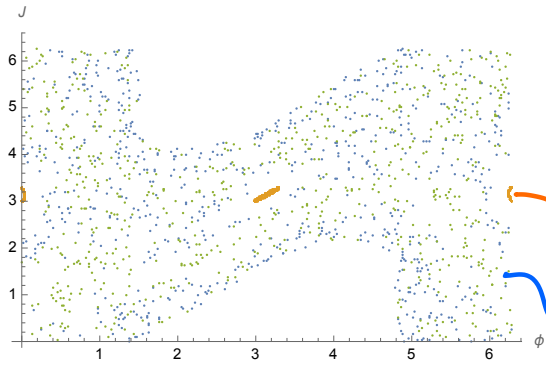
→ this seems "integrable"?



} which IC?? → chaos

"onset of chaos" at
 $\epsilon_c \sim 0.95$.

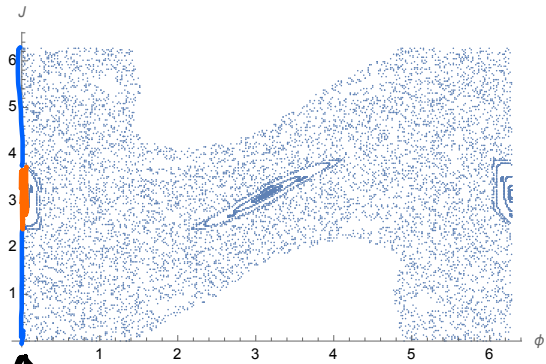
Numerics at $\epsilon = 1.8$.



→ "island of integrability"

→ blue / green ICs "mix"

chaos
integrable



↑ all dynamics started at $\phi_0 = 0$.