## Homework 12

Due: December 10 at 11:59 PM. Submit on Canvas.

Problem 1 (Bifurcations): Consider a "generic" one-dimensional map

$$x_{n+1} = f(x_n). \tag{1}$$

When the stability and/or number of fixed points of one of these maps changes abruptly as the function f is perturbed, we say that there is a **bifurcation**. In this problem, we will explore the kinds of generic bifurcations that can arise in one-dimensional maps.

- 15 A: Suppose that f(x) is an increasing function (i.e. f'(x) > 0) in the interval  $x_1 < x < x_2$ , and assume that  $f(x_1) \neq x_1$  and  $f(x_2) \neq x_2$ . Show that<sup>1</sup> if there are k unstable fixed points, there are either k-1, k or k+1 stable fixed points. Describe how, if you perturbed the function f(x) sufficiently weakly,<sup>2</sup> you could only ever create (or remove) a *pair* of two fixed points one stable, and one unstable.
- 15 B: Show that if f(x) is allowed to be decreasing as well, then it is possible for a stable fixed point to be perturbed into an unstable fixed point, without the creation of any new fixed point pair. By considering what would happen to the map  $f^{[2]}(x)$  from Lecture 38, argue that when a stable fixed point turns into an unstable fixed point, we will generically have a **period doubling bifurcation**, as in Lectures 38 and 39. Your result implies that the period doubling transition to chaos should be generic for one-dimensional maps.
- 15 C: As an explicit example of a perturbed map, consider the perturbed logistic map

$$x_{n+1} = rx_n(1 - x_n) + \epsilon \tag{2}$$

for 0 < r < 3.1 and  $\epsilon \ge 0$  sufficiently small. Calculate the location and stability of fixed points as a function of  $\epsilon$ , within this range of r. Show that when  $\epsilon = 0$ , the motion of fixed points under the perturbation of increasing r is unusual, and explain why. In contrast, show that for  $\epsilon > 0$ , the behavior of the map is "typical" and fits into the phenomenology you found above for "generic" maps.<sup>3</sup>

**Problem 2** (Quartic Hamiltonian): Consider canonical coordinates on phase space  $\mathbb{R}^4$ , with Hamiltonian

$$H = \frac{p_x^2 + p_y^2}{2} + \frac{x^4 + y^4 + 2\alpha x^2 y^2}{4}.$$
(3)

- 10 A: Take  $\alpha \ge 0$ . Where can you prove (easily!) the system is integrable?
- 15 B: Numerically look for integrability vs. chaos. Can you find compelling evidence for chaos at certain values of  $\alpha$ ? (You don't need to exhaustively characterize where there is chaos.)

<sup>&</sup>lt;sup>1</sup>*Hint:* Draw a graph of x and (any appropriate function) f(x) in the interval  $[x_1, x_2]$ . What are the possible choices of f(x) you need to consider?

<sup>&</sup>lt;sup>2</sup>*Hint:* The perturbed f(x) is still increasing.

<sup>&</sup>lt;sup>3</sup>Feel free to use numerical methods if necessary to discuss what happens to the period-doubling bifurcation near r = 3.

**Problem 3** (Solvable limit of chaos in the logistic map): Consider the logistic map of Lecture 37 at the maximal value r = 4:

$$x_{n+1} = 4x_n(1 - x_n). (4)$$

10 A: Suppose that we write

$$x_n = \frac{1 - \cos(2\pi y_n)}{2} \tag{5}$$

where we identify  $y_n \sim y_n + 1$  as a periodically identified variable. Plug (5) into (4) and use trigonometric identities to show that<sup>4</sup>

$$y_{n+1} = 2y_n \pmod{1} = 2y_n - \lfloor 2y_n \rfloor.$$
 (6)

10 B: Let us now show that (6) is an exactly solvable model of chaos. Without loss of generality, we may write

$$y_0 = \sum_{k=1}^{\infty} 2^{-k} b_k$$
 (7)

where each  $b_k = 0$  or 1. What is  $y_n$ ? Suppose that we start with two initial conditions  $y_0$  and  $y'_0$  obeying  $|y_0 - y'_0| = \epsilon \ll 1$ . At what step  $n \operatorname{can} |y_n - y'_n| \sim 1$ ? Compare the divergence of trajectories to the butterfly effect discussed in Lecture 40, and estimate a "Lyapunov exponent" of the logistic map at r = 4 analytically.

10 C: Calculate the box dimension (Lecture 41) of the set S, defined by

 $S = \{y_0 \text{ such that for all sufficiently large } n, y_n = 0\}.$ (8)

5 D: Now that you have seen the solution to the r = 4 logistic map, argue that for a typical trajectory not in S, the probability density for  $x_n$  (at large n, in the original x coordinates) is

$$\rho(x) = \frac{1}{\pi \sqrt{x(1-x)}}.$$
(9)

10 E: One way to define a "multifractal dimension"  $d_q$  is as follows. Consider  $\epsilon^{-1}$  boxes of size  $\epsilon$ , and define

$$\mu_j = \mathbb{P}[(j-1)\epsilon < x_n < j\epsilon] \tag{10}$$

to be the probability that x is found in box j. Define

$$d_q = \lim_{\epsilon \to 0} \frac{1}{1 - q} \frac{1}{\log \epsilon^{-1}} \sum_{i=1}^{\epsilon^{-1}} \mu_i^q.$$
 (11)

Calculate  $d_q$  for the r = 4 logistic map, and compare your answer to **C**. If you find that this definition of multifractal dimension  $d_q$  does detect the fractal structure of the logistic map, explain intuitively what is responsible for the effect.

<sup>&</sup>lt;sup>4</sup>Here a (mod 1) denotes only the "decimal" part of a, e.g.  $4.32 \pmod{1} = 0.32$ .