

PHYS 7810  
Hydrodynamics  
Spring 2024

Lecture 12  
Vortex dynamics

February 22

Today: dynamics of vorticity in Galilean-invariant fluids:

$$\omega_i = \epsilon_{ijk} \partial_j v_k \quad [\vec{\omega} = \nabla \times \vec{v}]$$

Assume flow incompressible:  $\nabla \cdot \vec{v} = 0 = \partial_i v_i$

Navier-Stokes equation: const. (assume)  $\leftarrow$  kinematic viscosity  

$$\partial_t v_i + v_j \partial_j v_i + \frac{1}{\rho} \partial_i P = \frac{\eta}{\rho} \partial_j \partial_j v_i = \nu \partial_j \partial_j v_i$$
 (air:  $\nu \sim 10^{-5} \frac{m^2}{s}$ )

$$\partial_t \omega_i + \epsilon_{ijk} \partial_j [v_l \partial_l v_k + \frac{1}{\rho} \partial_k P] = \nu \partial_l \partial_l \omega_i$$

$$\partial_t \omega_i + v_l \partial_l \omega_i + \underbrace{(\partial_l v_k)}_{\partial_k v_l + \epsilon_{lkm} \omega_m} \epsilon_{ijk} \partial_j v_l = \nu \partial_l \partial_l \omega_i$$

$$\partial_k v_l + \epsilon_{lkm} \omega_m$$

$$\begin{aligned} & \epsilon_{ijk} \cancel{(\partial_k v_l)} (\partial_j v_l) + \underbrace{\epsilon_{ijk} \epsilon_{lkm} \omega_m \partial_j v_l}_{-\omega_m \partial_m v_i + \cancel{\omega_i \partial_m v_m}} \end{aligned}$$

Conclude:

$$\partial_t \omega_i + v_l \partial_l \omega_i = \omega_m \partial_m v_i + \nu \partial_j \partial_j \omega_i$$

Limit of 2d flow:  $v_z = 0$  and  $\partial_z = 0$ .

$$\omega_z = \partial_x v_y - \partial_y v_x.$$

$$\partial_t \omega_z + \vec{v} \cdot \nabla \omega_z = \nu \nabla^2 \omega_z \quad [\text{but } \vec{v} \text{ depends on } \omega_z].$$

each vortex is connected

vorticity diffuses

Example 1: isolated "point vortex":

from lec 10:

$$v_x = \frac{\Gamma}{2\pi} \frac{-y}{x^2 + y^2}$$

$$v_y = \frac{\Gamma}{2\pi} \frac{x}{x^2 + y^2}$$

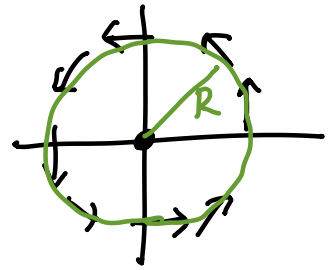
$$\hookrightarrow v_\theta = \frac{\Gamma}{2\pi r} \quad v_r = 0.$$

(polar coord)

$$\omega_z = \partial_x v_y - \partial_y v_x = \Gamma \delta(x) \delta(y)$$

$$\hookrightarrow \oint \vec{v} \cdot d\vec{s} = \frac{\Gamma}{2\pi R} \cdot 2\pi R = \Gamma$$

$$= \int \omega_z dx dy \quad \text{by Stokes' Thm.}$$



Incompressible:  $\nabla \cdot \vec{v} = 0$

$$\partial_t \omega_z + \vec{v} \cdot \nabla \omega_z = 0 \quad \int \frac{\partial \omega_z}{\partial \theta} = 0.$$

static point vortex is a solution.

Add viscosity?  $\partial_t \omega_z + \vec{v} \cdot \nabla \omega_z = \nu \nabla^2 \omega_z$   
rotational symmetry

$$\omega_z(r, \theta, t) = \frac{\Gamma}{4\pi \nu t} e^{-r^2/4\nu t} \quad (\text{cf lec. 2})$$

$$\hookrightarrow \text{Expect: } \vec{v} = v_\theta(r, t) \hat{\theta}$$

$$\nabla \times \vec{v} = \omega_z = \frac{1}{r} \partial_r (r v_\theta)$$

Integrate to get:  $\int_r^\infty dr' r' \omega_z(r') = r' v_\theta(r') \Big|_r^\infty$

$$\frac{\Gamma}{2\pi} e^{-r^2/4\nu t} = \frac{\Gamma}{2\pi} - r v_\theta(r)$$

So  $v_\theta(r) = \frac{\Gamma}{2\pi r} [1 - e^{-r^2/4\nu t}]$ .

Example 2: 2 <sup>point</sup> vortices, 2d...

Neglect dissipation:  $\partial_t \omega_z + \vec{v} \cdot \nabla \omega_z = 0$ .

Suppose  $\omega_z = \Gamma_1 \delta(\vec{r} - \vec{r}_1(t)) + \Gamma_2 \delta(\vec{r} - \vec{r}_2(t))$ .

$$-\Gamma_1 \dot{\vec{r}}_1 \cdot \nabla \delta(\vec{r} - \vec{r}_1) - \Gamma_2 \dot{\vec{r}}_2 \cdot \nabla \delta(\vec{r} - \vec{r}_2) + \Gamma_1 \vec{v}(\vec{r}_1) \cdot \nabla \delta(\vec{r} - \vec{r}_1) + \Gamma_2 \vec{v}(\vec{r}_2) \cdot \nabla \delta(\vec{r} - \vec{r}_2) = 0$$

$\hookrightarrow \dot{\vec{r}}_1 = \vec{v}(\vec{r}_1)$  and  $\dot{\vec{r}}_2 = \vec{v}(\vec{r}_2)$   
 contributions to  $\vec{v}$  from other vortices.

$$\frac{d\vec{r}_1}{dt} = \vec{v}_2(\vec{r}_1) = \frac{\Gamma_2}{2\pi} \frac{\hat{z} \times (\vec{r}_1 - \vec{r}_2)}{|\vec{r}_1 - \vec{r}_2|^2}$$

and  $\frac{d\vec{r}_2}{dt} = \frac{\Gamma_1}{2\pi} \frac{\hat{z} \times (\vec{r}_2 - \vec{r}_1)}{|\vec{r}_1 - \vec{r}_2|^2}$ .

Suppose  $\Gamma_1 = \Gamma_2 > 0$ :

Suppose  $\Gamma_1 = -\Gamma_2$



point vortices will move in a circle.



"vortex dipole" moves at const. velocity.

Generalize to  $n$  point vortices: vortex  $\alpha = 1, \dots, n$

$$\frac{d\vec{r}_\alpha}{dt} = \sum_{\beta (\neq \alpha)} \frac{\Gamma_\beta}{2\pi} \frac{\hat{z} \times (\vec{r}_\alpha - \vec{r}_\beta)}{|\vec{r}_\alpha - \vec{r}_\beta|^2}$$

Nice math problem!  $\rightarrow$  Hamiltonian mechanics!

$$H = - \sum_{\alpha \neq \beta} \frac{\Gamma_\alpha \Gamma_\beta}{2\pi} \log |\vec{r}_\alpha - \vec{r}_\beta| \quad \text{with} \quad \{x_\alpha, y_\beta\} = \Gamma_\alpha \delta_{\alpha\beta}$$

$\curvearrowright$   
canonical conj!

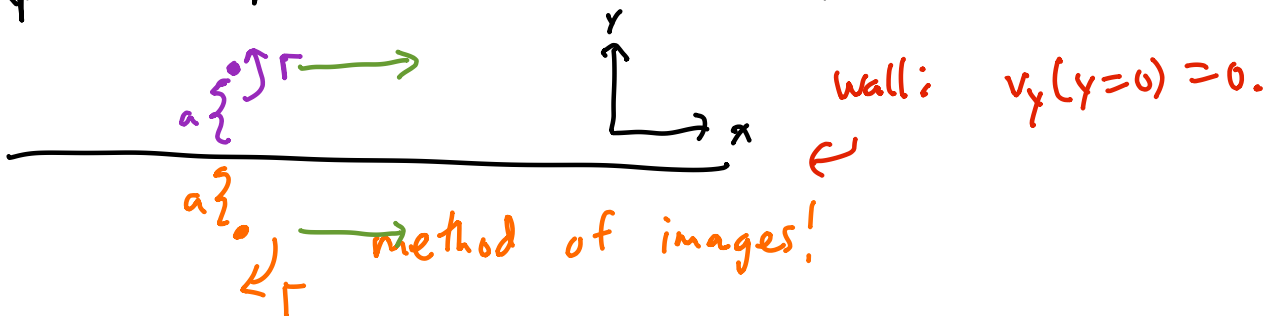
It is conserved... as is  $\sum \Gamma_\alpha x_\alpha, \sum \Gamma_\alpha y_\alpha, \sum \Gamma_\alpha (x_\alpha^2 + y_\alpha^2)$

Physics: • superfluid vortex dynamics

$\hookrightarrow$  quantized  $\Gamma = \pm \hbar/m$

• 2d charged particles in strong B-field (Lowest Landau level)

Example 3: point vortex near wall

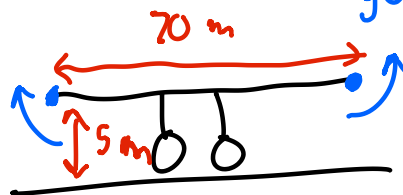
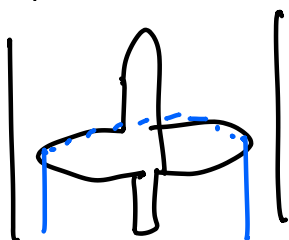


$$v_y = \frac{-\Gamma x}{x^2 + (y-a)^2} - \frac{-\Gamma x}{x^2 + (y+a)^2}$$

indeed:  $v_y(y=0) = 0$ .

Example 4: "wingtip vortex" of planes...

needed to generate lift!



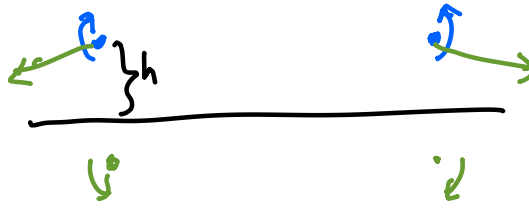
lead to dangerous flow for other planes.

→ viscous dissipation?

(air:  $\nu \sim 10^{-5} \text{ m}^2/\text{s}$ )

diffuse  $\sim 5\text{m} \sim \sqrt{\nu t} \rightarrow t \sim 2.5 \times 10^6 \text{ s}$

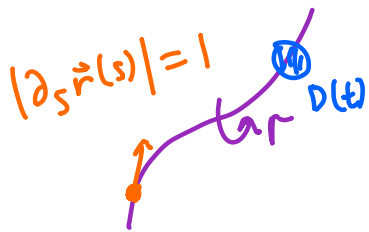
→ vortex motion:



estimate  $\Gamma \sim v_{\text{plane}} \cdot w \rightarrow 0.5\text{m}$  :  $v_{\text{vortex}} \sim \frac{\Gamma}{h} \sim 5 \text{ m/s}$  (too big)  
( $t \sim 1 \text{ s}$ )

Reality:  $\sim 1 \text{ min}$  (per FAA).

Now let's turn to 3d vortex line motion:



Analogous to Biot-Savart:

$$\vec{v}(\vec{x}) = \frac{\Gamma}{4\pi} \int ds \frac{\partial_s \vec{r} \times (\vec{x} - \vec{r}(s))}{|\vec{x} - \vec{r}(s)|^3}$$

parameterize by  $\vec{r}(s)$

Claim:  $\frac{\partial}{\partial t} \vec{r} = \vec{v}(\vec{r}, t)$ .

To justify:  $\frac{d}{dt} \oint_{\partial D(t)} \vec{v}(\vec{r}, t) \cdot d\vec{r} \stackrel{\text{Stokes}}{=} \frac{d}{dt} \int_D \vec{\omega} \cdot d\vec{A} \sim \frac{d}{dt} \Gamma$

$$= \oint_{\partial D} [\partial_t \vec{v} + \vec{v} \cdot \nabla \vec{v}] \cdot d\vec{r} + \oint_{\partial D} \vec{v} \cdot \frac{d\vec{r}}{dt}$$

$$= \oint_{\partial D} \left( -\frac{1}{\rho} \nabla P + \cancel{\nu \nabla^2 \vec{v}} \right) \cdot d\vec{r} + \oint_{\partial D} d\left(\frac{v^2}{2}\right)$$

ignore diss

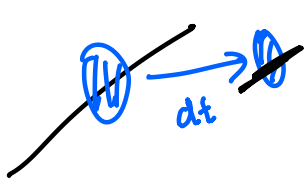
$$= 0 + 0 \quad \left( \oint d\vec{r} \cdot \nabla f = 0 \right)$$



Kelvin's Circulation Thm:

$$\frac{d}{dt} \int_{\partial D(t)} \vec{v} \cdot d\vec{s} = 0:$$

vortex line lasts forever...  
convects along w/ flow



circulation pierces this disk, so  
vortex line obeys

$$\frac{\partial \vec{r}}{\partial t} = \vec{v}(\vec{r}, t)$$