PHYS 7810 Hydrodynamics Spring 2024

Lecture 8

Navier-Stokes equations

February 8 Hydrody namics of liquids & gases dissipative EFT; degrees of freedom = conserved quant: (M) Particles / Mass / "charge" Lp) (Pi) Momentum (g_i) (६) (E) Energy (Lij) Angular Momentum (return in lec 9) total entropy Statistical stationary state? $P_{55} = e^{-\Phi} = e^{-BE} \rightarrow e^{S}$ e where $S = \int dx s(\varepsilon, p, g_i)$ count "field configs" available... d spatial dias. Thermodynamic review! First law of thermoin (kg=1) $T ds = d\epsilon - p_{H} d\rho - v_i dg_i \rightarrow \frac{\partial s}{\partial \epsilon} \Big|_{\rho,g} = \frac{1}{T}$

$$ds = \frac{1}{T} d\epsilon - \frac{\mu_{1}}{T} d\rho - \frac{v_{1}}{T} dg_{1}$$

$$Hermo conjugate: -v_{1}/T:$$

$$\mu_{5} = \frac{6\epsilon}{5g_{1}} = -\frac{5S}{5g_{2}} = \frac{v_{1}}{T}$$

$$Gilds = Dulem relation:$$

$$dE = TdS + \mu_{1}dM + v_{1}dP_{1} - PdV$$

$$work done by system + environ.$$

$$\frac{dE}{dV} = \frac{c}{c} = Ts + \mu_{1}p + v_{1}g_{1} - P$$

$$General MSR \ Lagrangian \ \mathcal{R}_{MSR}(\epsilon, p, g_{1}, \pi_{1}, \pi_{2}, \pi_{1})?$$

$$Time - reversal \ symmetry:$$

$$p \rightarrow p \qquad \epsilon \rightarrow \epsilon \qquad g_{1} \rightarrow -g_{1}$$

$$\pi_{p} \rightarrow -\pi_{p} + ipp \qquad \pi_{e} \rightarrow -\pi_{e} + i\mu_{e} \qquad \pi_{i} \rightarrow \pi_{i} - i\mu_{i}$$

$$\mathcal{L}_{MSR} = \pi_{p}\partial_{e}p + \pi_{e}\partial_{e}c + \pi_{i}\partial_{e}g_{1} - H_{MSR}(\cdots)$$

$$H_{MSR} = [\partial_{i}\pi_{p}]T^{i} + (\partial_{i}\pi_{e})E^{i} + (\partial_{i}\pi_{i})T_{i}]$$

$$due \ to \ conservation [laws of M_{i} \in P_{i}].$$

$$Stress \ tensor$$

$$J_{i} = -D\partial_{i}p + \cdots ?$$

Note:
$$\frac{S S_{NSR}}{S_{Tp}}\Big|_{\substack{N_{p} = 0}} = \partial_{t} p - \partial_{i} \left(-J^{i} \right) = \frac{1}{\partial_{t} p} + \frac{1}{\partial_{t}} J_{i} = 0$$

$$\int_{local conversation law}$$
First look for ideal hydro (dissipationless & $J_{i}(p, \epsilon, g)$
ind. of derivatives):
Time-reversal symmetry:
 $J_{i} \rightarrow -J_{i}$ $\mathcal{E}_{i} \rightarrow -\mathcal{E}_{i}$ $T_{ij} \rightarrow T_{ij}$
Assume rotation invariance...:
 $\mathcal{E}_{i} = -A(\epsilon, p, g; g; 1)v_{i}$; $J_{i} = B \cdot v_{i}$; $T_{ij} = C_{ij} + Dv_{i}v_{j}$
But impose $T - even:$
 $\mathcal{L}_{MSR} \rightarrow J_{MSR} + i \left[\mu_{e}\partial_{e}\epsilon + \mu_{p}\partial_{e}p + \mu_{i}\partial_{e}g_{i} - \mathcal{E}_{i}\partial_{i}\mu_{e} - J_{i}\partial_{i}\mu_{p} - T_{ij}\partial_{i}\mu_{j} \right]$
detailed balance, or thermo: $-\partial_{t}S$ $total divergence?
 $-\partial_{i}S^{i}$
 $= A_{i}v_{i}\partial_{i} + B_{i}v_{i}\partial_{i} + \frac{V_{i}v_{j}}{2H}$.
 $C is $C\left(\frac{1}{T}, \frac{\mu_{M}}{T}, \frac{v_{j}v_{j}}{2H}\right)$.
 $\partial_{i}C = \frac{\partial C}{\partial (V_{i})} = \partial_{i}(\frac{1}{T}) + \frac{\partial C}{\partial (mv_{i})} |_{v_{i}v_{i}v_{i}}\partial_{i} + \frac{\partial C}{\partial (\frac{v_{i}}{T})} |_{v_{i}v_{i}}\partial_{i} + \frac{\partial C}{\partial (\frac{$$$

 $0 = Av_{i} + \frac{2C}{2(4)} \left|_{v_{i}} v_{i}^{2} \cdot \frac{v_{i}}{T} - Dv_{i} \frac{v_{j}v_{j}}{T} + \frac{2C}{2(\frac{2}{3})} \right|_{v_{i}} v_{i} \frac{v_{j}v_{j}}{2T}$ $D = -B_{v_i} + \frac{v_i}{\tau} \frac{\partial C}{\partial (\mu_{f_i})} |_{V_{f_i} v^2}$ $0 = \frac{v_i v_j}{\tau} \frac{\partial C}{\partial (v_{h\tau})} \Big|_{T_{i,r}} - \frac{\partial v_i v_j}{\tau} \cdot \frac{v_{ehind}}{\lambda} \text{ of thermal}$ A, B, D fixed by one function C(E, p, g;). define to be pressure P. $D = \frac{1}{T} \frac{\partial P}{\partial (v_{T})} \Big|_{\frac{1}{T}} \frac{P}{\Phi} = \frac{\partial P}{\partial (v_{T})} \Big|_{\frac{T}{T}}$ Use Gibbs-Duhem & 1st law: $d(\epsilon + P) = d\epsilon \cdot dP = d(r_{h} p + v_{i}g_{i} + Ts)$ $\frac{\partial P}{\partial v_i} = g_i, \text{ then } \frac{\partial P}{\partial v_i} = v_i \frac{\partial P}{\partial (v_i)} \cdots g_i = D v_i \text{ (lec } q)$ Next: $B = \frac{1}{T} \frac{\partial P}{\partial (\frac{r}{T})} \Big|_{\frac{1}{T}, \frac{r}{T}} = \frac{\partial P}{\partial \mu_{h}} \Big|_{T, \frac{r}{T}} = \rho$ B=p Similar calculation: $A = \varepsilon + P$ Combine, constitutive relations for ideal fluid: $\tau_{ij} = PS_{ij} + Dv_iv_j$ $\mathcal{E}_{i} = (\varepsilon + P)v_{i}$ $J_{i} = \rho v_{i}$ $\partial_{\xi}\rho + \partial_{i}(\rho v_{i}) = 0$ $\partial_{\xi}g_{i} + \partial_{i}P + \partial_{j}(P v_{i}v_{j})$ $\partial_{\mu} \varepsilon + \partial_i ((\varepsilon + P) v_i) = 0$ Euler eq.

What are noise & dissipation addable to EFT?
- respect T
- add as few d; as possible. [ler 7: hydro and long wavelengths]
For momentum const:

$$2msk = \cdots + i \partial_i \pi_s T \eta_{ijkl} \partial_k (\pi_l - i \eta_l)$$

 $T_{iscosily} tensor : rotation invariance:$
 $\eta_{ijkl} = \eta (f_{ik} f_{ijk} + f_{ilk} f_{ijk}) - \frac{1}{2} f_{ij} f_{kl})$
 $for energy & mass/charge:
 $2msk = \cdots + i T \partial_i (\pi_k \pi_p) \left(A_{kp}^{ss} A_{kp}^{cp} \right) \partial_i (\pi_k - i \eta_p)$
 $incoherest conductivity$
constitutive relations:
 $T_i = pv_i - T A_{kp}^{pr} \partial_i (\frac{f_n}{T}) - T A_{kp}^{cp} \partial_i (\frac{1}{T})$
 $= pv_i - T A_{kp}^{pr} \partial_i (\frac{f_n}{T}) - T A_{kp}^{cp} \partial_i (\frac{1}{T})$
 $E_i = (\epsilon + P)v_i - T a_{ijkr} - \bar{k} \partial_i T$
 $V_{ijkl} \partial_k (T)$
Navier-Stokes:
 $\partial_k p + \partial_i T = 0$ $2 + \epsilon + 2i \mathcal{E}_i = 0$ $2 + g_i + \partial_i \tau_i = 0$.
 $D_{ur} = EFT$ gives us $2nd$ hav of thermodynamics:
 $\partial_k s + \partial_i S^i \ge 0$.
 $i.e. \frac{d_i}{d_i} d^a x s \ge 0$.$