$electromagnetism \rightarrow electromagnetic radiation$ 

## Transition Radiation

A charge q moves at constant velocity  $\mathbf{v} = v\hat{\mathbf{x}}$  from a vacuum (x < 0) into a linear dielectric (x > 0) with index of refraction  $n = \sqrt{\epsilon/\epsilon_0}$ , which you can assume is real. Assume that v < c/n.



(a) Assume that at time t = 0, the charge is at x = y = z = 0. Fourier transform Maxwell's equations in the y, z and t directions, and show that the electric field  $\mathbf{E}_{s}$  generated by the moving source, ignoring the presence of the boundary, is given by

$$\mathbf{E}_{s}(x,k_{y},k_{z},\omega) = \frac{\mathrm{i}q}{\epsilon_{0}n^{2}v} \frac{\omega v \left(\frac{n^{2}}{c^{2}} - \frac{1}{v^{2}}\right) \hat{\mathbf{x}} - \mathbf{k}}{k^{2} - \omega^{2} \left(\frac{n^{2}}{c^{2}} - \frac{1}{\omega^{2}}\right)} \mathrm{e}^{\mathrm{i}\omega x/v}$$

In this equation, take  $n^2$  to be a function of x.

(b) As you may have noticed, the above does not quite solve Maxwell's equations, because of the boundary at x = 0, where the abrupt change in n will cause  $\mathbf{E}_s$  to be discontinuous. We can remedy this by writing the true electric field as  $\mathbf{E} = \mathbf{E}_s + \mathbf{E}_r$ , where  $\mathbf{E}_r$  is a radiation field, added to satisfy the boundary conditions. Show, on general principles, that the form of  $\mathbf{E}_r$  must be

$$\mathbf{E}_{\mathbf{r}} = (x, k_y, k_z, \omega) = \mathrm{i}A_{\pm}(\mathbf{k}, \omega) \left[ \hat{\mathbf{x}} \mp \frac{\mathbf{k}}{k^2} \sqrt{\frac{n^2 \omega^2}{c^2} - k^2} \right] \mathrm{e}^{\pm \mathrm{i}x \sqrt{n^2 \omega^2 / c^2 - k^2}}$$

where the top sign is used for x > 0, and the bottom sign for x < 0.

(c) The boundary condition at x = 0 provides equations which can be solved to find  $A_{\pm}$ . Write down these equations, but do not solve them.

If you did the nasty algebra, you would find

$$A_{-} = \frac{q\beta\kappa^{2}(n^{2}-1)(1-\beta^{2}+\beta\sqrt{n^{2}-\kappa^{2}})}{\omega(1-\beta^{2}+\beta^{2}\kappa^{2})(1+\beta\sqrt{n^{2}-\kappa^{2}})(\sqrt{n^{2}-\kappa^{2}}+n^{2}\sqrt{1-\kappa^{2}})}$$

where  $\beta \equiv v/c$  and  $\kappa \equiv kc/\omega$ .

For the remainder of the problem, our goal will be to compute the energy W radiated by this particle as it moves across the boundary, in the limit of  $t \to \infty$ . This is the so called **transition radiation**. To do this, we begin with the formula (note that the factor of 2 comes from the contribution of the magnetic field):

$$W = \int_{x<0} \mathrm{d}V \ \epsilon_0 \mathbf{E}_{\mathbf{r}} \cdot \overline{\mathbf{E}_{\mathbf{r}}}.$$

This is basically because in the large time limit, only the radiation field is present at x < 0. You do not need to justify this statement, although it is not difficult.

(d) Begin by showing that

$$W = \int d^2 \mathbf{k} d\omega \ \frac{\epsilon_0}{(2\pi)^4} \frac{\omega^2}{ck^2} \sqrt{1 - \frac{c^2 k^2}{\omega^2}} |A_-|^2.$$

(e) Now justify, the following statement: the only relevant integral is over momenta such that  $ck \leq \omega$ . Switching variables to

$$\sin\theta \equiv \frac{ck}{\omega},$$

re-write the integral for W as an integral only over  $\theta$  and  $\omega$ .

(f) Work in the limit where  $n \approx 1$  and  $v \approx c$ . Show that, in this limit, we can approximate

$$W \sim \left(\frac{n-1}{n+1}\right)^2 \log \frac{1}{1-v^2/c^2}.$$

Comment on the result. How much of the particle's energy (relatively) is lost to transition radiation as it crosses the boundary, as the particle becomes more energetic?