classical mechanics  $\rightarrow$  rigid body motion

## Motion on SO(d) by Lagrange Multipliers

Previously, we have seen how to derive the equations of motion for the free *d*-dimensional rigid body by appealing to the group structure of the configuration space SO(d). There is a less mathematically elegant, but more direct approach, which employs Lagrange multipliers.

As before, let us consider the motion of a rigid body about the origin. If A describes the matrix in the configuration space SO(d), we have previously explained why the Lagrangian is

$$L = \operatorname{tr}\left[\frac{1}{2}\dot{A}K\dot{A}^{\mathrm{T}} - B\left(A^{\mathrm{T}}A - 1\right)\right]$$

where

$$K \equiv \int \mathrm{d}^d x \; x x^\mathrm{T} \rho(x).$$

Of course, the Lagrange multiplier B leads to a trivial term if we choose to express A in terms of the Lie algebra-valued velocities, but in this problem we will not do this. Instead, we will allow A to be completely arbitrary, and only impose the constraint that  $A \in SO(d)$  by the Lagrange multiplier (and continuity of the equation of motion).

- (a) Explain why B is a symmetric matrix.
- (b) Show that the equations of motion imply

$$A^{\mathrm{T}}\ddot{A}K - K\ddot{A}^{\mathrm{T}}A = 0.$$

(c) Define the matrix

$$\Omega \equiv A^{\mathrm{T}} \dot{A}.$$

Verify that  $\Omega$  is antisymmetric.

- (d) What is the equation of motion in terms of  $\Omega$ ?
- (e) Verify by further manipulations that you can recover the equation we derived by group theoretic arguments earlier:

$$M^{ab}\dot{\omega}^b = f^{adc}M^{db}\omega^b\omega^c$$

where  $\omega^b$  are the generalized angular velocities,  $M^{ab}$  is the generalized moment of inertia, and  $f^{adc}$  are the structure constants of SO(d).